

Dimensional reduction on curved manifolds: low energy, swampland, de Sitter

David ANDRIOT

CERN, Geneva, Switzerland

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[arXiv:1806.10999](#)

Work in progress

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Motivation: **connection of string theory to pheno.**
(cosmology, particle physics, gravitational waves)
Context: string compactifications: $10d \rightarrow 4d$

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1. Find a 10d string background:

4d space-time (de Sitter, Minkowski) \times 6d compact manifold \mathcal{M}

2.

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Here: classical (perturbative) backgrounds:
solutions of 10d type IIA/B supergravities with
 D_p -branes/orientifolds O_p -planes

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2. Dimensional reduction:

fluctuations around background

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fluctuations around background, infinite (towers) of them

\leftrightarrow truncation to a finite set

$\rightarrow \int_{6d} \rightarrow$ 4d theory \leftrightarrow observations?

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Truncation + 4d theory: **low energy??**

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Consider 6d \mathcal{M} which are **curved/non-Ricci flat**.

Away from the lamppost, different regions of the landscape,
new appealing phenomenological features:

- stabilize some specific moduli
- (some classical) de Sitter solutions require $\mathcal{R}_6 < 0$

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 \rightsquigarrow **different behaviors/counterexamples w.r.t.
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Prime example: \mathcal{M} being a **group manifold**:
most are curved, easy to handle.

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Prime example: \mathcal{M} being a **group manifold**:

most are curved, easy to handle.

\mathcal{M} = Lie group / lattice , compact.

Underlying Lie algebra:

$$[V_a, V_b] = f^c{}_{ab} V_c \leftrightarrow de^a = -\frac{1}{2} f^a{}_{bc} e^b \wedge e^c$$

$$\hookrightarrow \mathcal{M}: ds_6^2 = \delta_{ab} e^a e^b, \quad e^a = e^a{}_m dx^m, \quad f^a{}_{bc} = 2\omega_{[b}{}^a{}_{c]}.$$

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Known de Sitter, Minkowski, anti-de Sitter solutions with \mathcal{M} =group manifold (and fluxes, D_p/O_p).

\Rightarrow dimensional reduction \rightarrow fluctuations + truncation ?

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Plan of the talk:

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Plan of the talk:

1. Fluctuations: $\Phi(x^\mu, x^m) = \sum_I \phi^I(x^\mu) U_I(x^m)$
 \leftrightarrow find an appropriate basis $\{U_I\}$

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1. Fluctuations: $\Phi(x^\mu, x^m) = \sum_I \phi^I(x^\mu) U_I(x^m)$
 \leftrightarrow find an appropriate basis $\{U_I\}$: Laplacian $\Delta_{\mathcal{M}}$ eigenmodes
 \Rightarrow spectrum of Laplacian on \mathcal{M} = nilmanifold.

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3. Recent de Sitter swampland criterion.

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Nilmanifold geometry

3d Heisenberg Lie algebra:

$$[V_1, V_2] = -\mathbf{f} V_3, \quad [V_1, V_3] = [V_2, V_3] = 0, \quad \mathbf{f} = -f^3_{12}$$
$$de^3 = \mathbf{f} e^1 \wedge e^2; \quad de^1 = 0; \quad de^2 = 0$$

\hookrightarrow 3d nilmanifold M .

Build 6d $\mathcal{M} = M \times M', M \times T^3$.

Coordinates with (positive) radii $r^{m=1,2,3}$:

$$e^1 = r^1 dx^1; \quad e^2 = r^2 dx^2; \quad e^3 = r^3 (dx^3 + N x^1 dx^2)$$

$$N = \frac{r^1 r^2}{r^3} \mathbf{f} \in \mathbb{Z}^*$$

$$V = \int d^3 x \sqrt{g} = r^1 r^2 r^3$$

Geometry: **fibration of a S^1_3 on a base $T^2_{1,2}$.**

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Eigenmodes		Eigenvalues
Scalars	$v_{p,q}$	$-\mu_{p,q}^2$
	$u_{k,l,n}$	$-M_{k,l,n}^2$
Exact one-forms	$dv_{p,q} (pq \neq 0)$	$-\mu_{p,q}^2$
	$du_{k,l,n}$	$-M_{k,l,n}^2$
Co-closed one-forms	$B_{\pm}^{p,q}$	$Y_{\pm}^{p,q}$
	$B_{\pm}^{k,l,n}$	$Y_{\pm}^{k,l,n}$

2- and 3-forms obtained by Hodge star on 1-forms and scalars
 \Rightarrow **complete spectrum of Δ_M** .

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2- and 3-forms obtained by Hodge star on 1-forms and scalars
 \Rightarrow **complete spectrum of Δ_M** .

$$\begin{aligned}(\Delta + \mu_{p,q}^2)v_{p,q} &= 0, & (\Delta + M_{k,l,n}^2)u_{k,l,n} &= 0 \\ \Delta dv_{p,q} &= -\mu_{p,q}^2 dv_{p,q}, & \Delta du_{k,l,n} &= -M_{k,l,n}^2 du_{k,l,n}\end{aligned}$$

$$v_{p,q}(x^1, x^2)$$

$$u_{k,l,n}(x^1, x^2, x^3)$$

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$$\Delta dv_{p,q} = -\mu_{p,q}^2 dv_{p,q}, \quad \Delta du_{k,l,n} = -M_{k,l,n}^2 du_{k,l,n}$$

$$v_{p,q}(\mathbf{x}^1, \mathbf{x}^2) = \frac{1}{\sqrt{V}} e^{2\pi i p x^1} e^{2\pi i q x^2}, \quad p, q \in \mathbb{Z}$$

$$\mu_{p,q}^2 = p^2 \left(\frac{2\pi}{r^1}\right)^2 + q^2 \left(\frac{2\pi}{r^2}\right)^2$$

$$u_{k,l,n}(\mathbf{x}^1, \mathbf{x}^2, \mathbf{x}^3)$$

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$$\mu_{p,q}^2 = p^2 \left(\frac{2\pi}{r^1}\right)^2 + q^2 \left(\frac{2\pi}{r^2}\right)^2, \quad M_{k,l,n}^2 = k^2 \left(\frac{2\pi}{r^3}\right)^2 + (2n+1)|k| \frac{2\pi|f|}{r^3}$$

$$u_{k,l,n}(\mathbf{x}^1, \mathbf{x}^2, \mathbf{x}^3) =$$

$$\sqrt{\frac{r^2}{|N|V}} \frac{1}{\sqrt{2^n n! \sqrt{\pi}}} e^{2\pi i k(x^3 + N x^1 x^2)} e^{2\pi i l x^1} \sum_{m \in \mathbb{Z}} e^{2\pi i k m x^1} \Phi_n^\lambda(w_m)$$

where $\Phi_n^\lambda(z) = |\lambda|^{\frac{1}{4}} e^{-\frac{1}{2}|\lambda|z^2} H_n(|\lambda|^{\frac{1}{2}}z)$, $k \in \mathbb{Z}^*$, $n \in \mathbb{N}$

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$$\Delta B = YB$$

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$$\Delta B = YB$$

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$$p^2 + q^2 \neq 0 : \quad \mathbf{B}_{\pm}^{\mathbf{p},\mathbf{q}} = D v_{p,q} \left(Qe^1 - Pe^2 + \frac{\Upsilon_{\pm}^{\mathbf{p},\mathbf{q}} - (P^2 + Q^2)}{\text{if}} e^3 \right)$$

$$p = q = 0 : \quad \begin{aligned} B_1^{0,0} &= v_{0,0} e^1, & B_2^{0,0} &= v_{0,0} e^2 & (Y_-^{0,0} &= 0) \\ B_3^{0,0} &= v_{0,0} e^3 & (Y_+^{0,0} &= \mathbf{f}^2) \end{aligned}$$

$$(P = p \frac{2\pi}{r_1}, \quad Q = q \frac{2\pi}{r_2})$$

$$\mathbf{Y}_{\pm}^{\mathbf{p},\mathbf{q}} = P^2 + Q^2 + \frac{\mathbf{f}^2}{2} \pm \sqrt{\left(P^2 + Q^2 + \frac{\mathbf{f}^2}{2}\right)^2 - (P^2 + Q^2)^2}$$

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$$B_3^{0,0} = v_{0,0} e^3 \quad (Y_+^{0,0} = \mathbf{f}^2)$$

$$(P = p \frac{2\pi}{r^1}, \quad Q = q \frac{2\pi}{r^2})$$

$$\mathbf{Y}_{\pm}^{\mathbf{p},\mathbf{q}} = P^2 + Q^2 + \frac{\mathbf{f}^2}{2} \pm \sqrt{\left(P^2 + Q^2 + \frac{\mathbf{f}^2}{2}\right)^2 - (P^2 + Q^2)^2}$$

$$\mathbf{B}_{\pm}^{\mathbf{k},\mathbf{l},\mathbf{n}} = \sum_{a=1}^3 \varphi_a^{\mathbf{k},\mathbf{l},\mathbf{n}}(x^m) e^a$$

$$\varphi_1^{\mathbf{k},\mathbf{l},\mathbf{n}} = \frac{1}{2} c \sqrt{2^n n!} \left(u_{\mathbf{k},\mathbf{l},\mathbf{n}} - \frac{\sqrt{(n+1)(n+2)} \mathbf{f}^2}{\alpha_n + 2|\lambda| + (n+2) \mathbf{f}^2} u_{\mathbf{k},\mathbf{l},\mathbf{n}+2} \right)$$

$$\varphi_2^{\mathbf{k},\mathbf{l},\mathbf{n}} = -\frac{i}{2} \text{sgn}(\lambda) c \sqrt{2^n n!} \left(u_{\mathbf{k},\mathbf{l},\mathbf{n}} + \frac{\sqrt{(n+1)(n+2)} \mathbf{f}^2}{\alpha_n + 2|\lambda| + (n+2) \mathbf{f}^2} u_{\mathbf{k},\mathbf{l},\mathbf{n}+2} \right)$$

$$\varphi_3^{\mathbf{k},\mathbf{l},\mathbf{n}} = -\frac{1}{2|\lambda|^{\frac{1}{2}}} c \sqrt{2^n n!} \frac{\mathbf{f} \sqrt{2(n+1)(\alpha_n + 2|\lambda|)}}{\alpha_n + 2|\lambda| + (n+2) \mathbf{f}^2} u_{\mathbf{k},\mathbf{l},\mathbf{n}+1},$$

$$\mathbf{Y}_{\pm}^{\mathbf{k},\mathbf{l},\mathbf{n}} = k^2 \left(\frac{2\pi}{r^3} \right)^2 + (2n+3) \frac{2\pi}{r^3} |k\mathbf{f}| + \frac{1}{2} \mathbf{f}^2$$

$$\pm \sqrt{\left(\frac{2\pi}{r^3} |k\mathbf{f}| + \frac{1}{2} \mathbf{f}^2 \right)^2 + 2(n+1) \frac{2\pi}{r^3} |k\mathbf{f}| \mathbf{f}^2}$$

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Truncations and the swampland

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Consistent truncations and gauged supergravity

Standard truncation on group manifolds:

Scherk–Schwarz truncation:

keep the left-invariant/Maurer–Cartan forms: $\text{constant} \times e^a$.

(compare to previous spectrum: $v_{p,q} \times e^{1,2,3}$, $v_{0,0} \times e^{1,2,3}$)

\hookrightarrow 4d theory: **gauged supergravity**, gaugings: $f^a{}_{bc}$.

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\hookrightarrow 4d theory: **gauged supergravity**, gaugings: $f^a{}_{bc}$.

Example with an $SU(3)$ structure:

keep finite set of forms on \mathcal{M} : $1, \omega_I, \alpha_K, \beta^K, \tilde{\omega}^I, \text{vol}_6$, on which fields are developed

+ build $SU(3)$ structure forms J, Ω , analogous to CY

$\hookrightarrow \mathcal{N} = 2$ gauged supergravity.

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Scherk–Schwarz truncations are **consistent truncations**:

select set of “independent” or decoupled modes \rightarrow 4d theory.

However: **not low energy** truncation/not low energy theory:

light modes are truncated, heavy modes are kept.

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Low energy truncation on the nilmanifold

$$-f^3_{12} = \mathbf{f} = \frac{r^3 N}{r^1 r^2}$$

\Leftrightarrow consider approximation: $r^3 \leq |N| r^3 \ll r^1, r^2$

(small fiber/large base)

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(**small fiber/large base**)

\Rightarrow **hierarchy** of scales: $|\mathbf{f}| \ll \frac{1}{r^1}, \frac{1}{r^2} \ll \frac{1}{r^3}$

Hierarchy between curvature vs radii scales,
geometric flux vs Kaluza–Klein scales.

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Low energy approximation: truncate masses $\geq \frac{1}{r^1}, \frac{1}{r^2}, \frac{1}{\sqrt{r^1 r^2}}$

\rightarrow apply on the complete spectrum of Laplacian

Low energy truncation on the nilmanifold

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Low energy approximation: truncate masses $\geq \frac{1}{r^1}, \frac{1}{r^2}, \frac{1}{\sqrt{r^1 r^2}}$

\rightarrow apply on the complete spectrum of Laplacian,
left with light modes (mass = 0 or $|\mathbf{f}|$):

$$1, \quad e^1, e^2, e^3, \quad e^2 \wedge e^3, e^3 \wedge e^1, e^1 \wedge e^2, \quad e^1 \wedge e^2 \wedge e^3$$

up to normalisation constant $v_{0,0} = \frac{1}{\sqrt{V}}$.

Low energy truncation on the nilmanifold

$$-f^3_{12} = \mathbf{f} = \frac{r^3 N}{r^1 r^2}$$

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up to normalisation constant $v_{0,0} = \frac{1}{\sqrt{V}}$.

Surprise: low energy truncation matches Scherk–Schwarz trunc.
(on the Laplacian spectrum, in this regime)

Specific to nilmanifold w.r.t. to other group manifolds.

Testing the (refined) swampland distance conjecture

In field space, move from ϕ_0 to $\phi_0 + \Delta\phi$:
a tower of modes of mass $m(\phi)$ becomes light:

$$m(\phi_0 + \Delta\phi) = m(\phi_0) f(\phi_0, \Delta\phi) e^{-\alpha \frac{\Delta\phi}{M_P}}$$

If at ϕ_0 : \checkmark effective theory \Rightarrow ruined at $\phi_0 + \Delta\phi$.

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Testing the (refined) swampland distance conjecture

In field space, move from ϕ_0 to $\phi_0 + \Delta\phi$:
a tower of modes of mass $m(\phi)$ becomes light:

$$m(\phi_0 + \Delta\phi) = m(\phi_0) f(\phi_0, \Delta\phi) e^{-\alpha \frac{\Delta\phi}{M_p}}$$

If at ϕ_0 : \checkmark effective theory \Rightarrow ruined at $\phi_0 + \Delta\phi$.

Test this here on towers of Laplacian eigenmodes:

points in field space: $r^3 \ll r^{1,2}$ $\xleftrightarrow{\text{dist.}}$ \bullet $\xleftrightarrow{\text{dist.}}$ \bullet

number of light modes: finite

masses of light modes: $0, |\mathbf{f}|$

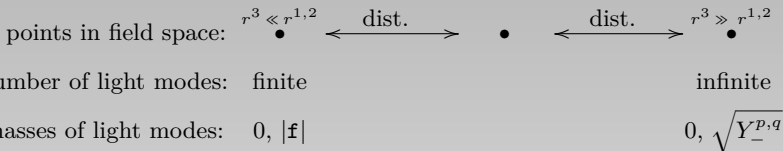
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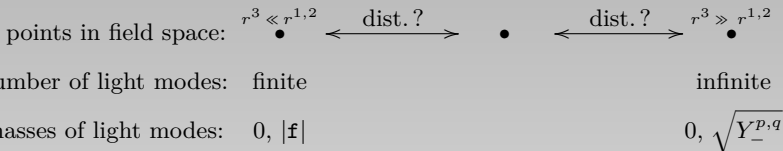
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Counterexample of the conjecture?

String effects? Quantum effects?

Due to the nilmanifold geometry, away from the lamppost...

De Sitter swampland criterion

Consider gravity minimally coupled to scalar fields with potential $V(\phi_i)$:

$$\mathcal{S} = \int d^4x \sqrt{|g_4|} (\mathcal{R}_4 + \text{kin. terms} - V)$$

Solutions with constant scalars:

$$\mathcal{R}_{\mu\nu} - \frac{g_{\mu\nu}}{2} \mathcal{R}_4 = -\frac{g_{\mu\nu}}{2} V \Rightarrow \mathcal{R}_4 = 2V, \quad \partial_{\phi_i} V = 0$$

Extrema of potential, value $V|_0$: maximally symmetric 4d space-time, cosmological constant $\Lambda = \frac{1}{2}V|_0$, de Sitter $V|_0 > 0$.

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Recent **swampland criterion**: any low energy effective theory of a consistent quantum gravity (i.e. not in the swampland) should verify the criterion

$$cV \leq |\nabla V|$$

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\Rightarrow extremum: $V|_0 \leq 0 \Rightarrow$ **no de Sitter solution**
(+ cosmological consequences)

Inspired by situation of de Sitter vacua in string theory:
difficult/unnatural.

Searches have provided similar conditions/criteria.

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Distinction: extremum/solution: $\partial_{\phi_i} V = 0$ and
vacuum/(meta)stable solution: $+\partial_{\phi_i}^2 V > 0$.

Difficulty is to have $V > 0$, $\partial_{\phi_i} V = 0$, $\partial_{\phi_i}^2 V > 0$ **together**

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Difficulty is to have $V > 0$, $\partial_{\phi_i} V = 0$, $\partial_{\phi_i}^2 V > 0$ **together**
 \leftrightarrow a natural criterion:

$\exists b_i \in \mathbb{R}, c_i \in \mathbb{R}_+$ such that

$$V + \sum_i b_i \phi_i \partial_{\phi_i} V + \sum_i c_i \phi_i^2 \partial_{\phi_i}^2 V \leq 0$$

$$\text{Solution: } V|_0 + \sum_i c_i (\phi_i^2 \partial_{\phi_i}^2 V)|_0 \leq 0$$

\Rightarrow **no stable de Sitter solution**, tachyonic de Sitter sol. \checkmark .

\Rightarrow checks? Cosmological implications?

There **exist** (unstable) 10d classical de Sitter **solutions**:
in type IIA/B with intersecting D_p/O_p , on group manifolds

$$\Rightarrow |\nabla V| \geq cV \text{ wrong ?}$$

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Argument saying that **two points can be compatible**:
Known 10d de Sitter solutions + 4d low energy effective theory
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↔ which de Sitter swampland criterion is valid? (if any)

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David
ANDRIOT

- Obtained Laplacian spectrum on a nilmanifold
- Found a low energy truncation that matches Scherk–Schwarz truncation; probably only for nilmanifolds
- Used the spectrum to “test” (refined) swampland distance conjecture, maybe a counterexample
- Proposed a new de Sitter swampland criterion that includes the notion of stability
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Thank you for your attention!

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