

Chiral Global Embedding of Fibre Inflation

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StringPheno 2018



ALMA MATER STUDIORUM
UNIVERSITÀ DI BOLOGNA



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[1709.01518] in collaboration with: M. Cicoli, D. Ciupke, V. Diaz, F. Muia and P. Shukla

Aim of the project

- Embed FI into consistent Calabi Yau orientifolds with D-branes and fluxes
- Get full moduli stabilisation
- Realise a chiral visible sector

Requirements:

- $h_{1,1} = 4$, Volume form $\mathcal{V} = \alpha_1 \sqrt{\tau_1 \tau_2 \tau_3} - \alpha_2 \tau_s^{3/2}$ [Cicoli,Muia,Shukla]
- **Orientifold involution** + $D3/D7$ -**brane setup** + tadpole cancellation
- World volume fluxes turned on, **chiral matter** on $D7$
- Freed-Witten anomaly cancellation + just one FI term
- No chiral intersection between visible sector and del Pezzo divisor
- **Moduli stabilisation** and inflation inside the Kähler cone
- **EFT** under control (Hierarchy of scales)

	x_1	x_2	x_3	x_4	x_5	x_6	x_7	x_8
4	0	0	0	1	1	0	0	2
4	0	0	1	0	0	1	0	2
4	0	1	0	0	0	0	1	2
8	1	0	0	1	0	1	1	4
	dP ₇	NdP ₁₁	NdP ₁₁	K3	NdP ₁₁	K3	K3	SD

- $SR1 = \{x_1x_4, x_1x_6, x_1x_7, x_2x_7, x_3x_6, x_4x_5x_8, x_2x_3x_5x_8\}$
- Intersection polynomial : $I_3 = 2D_4D_6D_7 - 2D_1^3$
- Volume form: $\mathcal{V} = 2t_4t_6t_7 + \frac{t_1^3}{3} = \frac{\sqrt{\tau_4\tau_6\tau_7}}{\sqrt{2}} - \frac{\tau_1^{3/2}}{3}$
- Kähler cone: $t_1 < 0$; $t_1 + t_7 > 0$; $t_1 + t_4 > 0$; $t_1 + t_6 > 0$
- Second Chern class: $c_2(X) = D_4D_5 + 4D_5^2 + 12D_5D_6 + 12D_5D_7 + 12D_6D_7$
- Topological quantities: $\Pi_i = \int_X c_2 \wedge \hat{D}_i$

Intersection curves and volumes:

	D_1	D_2	D_3	D_4	D_5	D_6	D_7	D_8
D_1	C_3	T^2	T^2	\emptyset	T^2	\emptyset	\emptyset	C_3
D_2	T^2	$P^1 \sqcup P^1$	$P^1 \sqcup P^1$	T^2	$P^1 \sqcup P^1$	T^2	\emptyset	C_3
D_3	T^2	$P^1 \sqcup P^1$	$P^1 \sqcup P^1$	T^2	$P^1 \sqcup P^1$	\emptyset	T^2	C_3
D_4	\emptyset	T^2	T^2	\emptyset	\emptyset	T^2	T^2	C_9
D_5	T^2	$P^1 \sqcup P^1$	$P^1 \sqcup P^1$	\emptyset	$P^1 \sqcup P^1$	T^2	T^2	C_3
D_6	\emptyset	T^2	\emptyset	T^2	T^2	\emptyset	T^2	C_9
D_7	\emptyset	\emptyset	T^2	T^2	T^2	T^2	\emptyset	C_9
D_8	C_3	C_3	C_3	C_9	C_3	C_9	C_9	C_{81}

	D_1	D_2	D_3	D_4	D_5	D_6	D_7	D_8
D_1	$2t_1$	$-2t_1$	$-2t_1$	0	$-2t_1$	0	0	$-4t_1$
D_2	$-2t_1$	$2t_1$	$2(t_1 + t_4)$	$2t_6$	$2(t_1 + t_6)$	$2t_4$	0	$4(t_1 + t_4 + t_6)$
D_3	$-2t_1$	$2(t_1 + t_4)$	$2t_1$	$2t_7$	$2(t_1 + t_7)$	0	$2t_4$	$4(t_1 + t_4 + t_7)$
D_4	0	$2t_6$	$2t_7$	0	0	$2t_7$	$2t_6$	$4(t_6 + t_7)$
D_5	$-2t_1$	$2(t_1 + t_6)$	$4(t_1 + t_7)$	0	$2t_1$	$2t_7$	$2t_6$	$4(t_1 + t_6 + t_7)$
D_6	0	$2t_4$	0	$2t_7$	$2t_7$	0	$2t_4$	$4(t_4 + t_7)$
D_7	0	0	$2t_4$	$2t_6$	$2t_6$	$2t_4$	0	$4(t_4 + t_6)$
D_8	$-4t_1$	$4(t_1 + t_4 + t_6)$	$4(t_1 + t_4 + t_7)$	$4(t_6 + t_7)$	$4(t_1 + t_6 + t_7)$	$4(t_4 + t_7)$	$4(t_4 + t_6)$	$8(t_1 + 2(t_4 + t_6 + t_7))$

Orientifold involution:

σ	O7	O3	N_{O3}	$\chi(O7)$	χ_{eff}
$x_1 \rightarrow -x_1$	D_1	$\{D_2D_3D_4, D_2D_4D_6, D_2D_5D_6, D_3D_4D_7, D_3D_5D_7, D_4D_6D_7, D_5D_6D_7\}$	14	10	-184
$x_2 \rightarrow -x_2$	$D_2 \sqcup D_7$	$D_1D_3D_5$	2	38	-192
$x_2 \rightarrow -x_3$	$D_3 \sqcup D_6$	$D_1D_2D_5$	2	38	-192
$x_4 \rightarrow -x_4$	$D_4 \sqcup D_5$	$D_1D_2D_3$	2	38	-192
$x_5 \rightarrow -x_5$	$D_4 \sqcup D_5$	$D_1D_2D_3$	2	38	-192
$x_6 \rightarrow -x_6$	$D_3 \sqcup D_6$	$D_1D_2D_5$	2	38	-192
$x_7 \rightarrow -x_7$	$D_2 \sqcup D_7$	$D_1D_3D_5$	2	38	-192
$x_8 \rightarrow -x_8$	D_8	\emptyset	0	208	-28

• Focus on: $x_8 \rightarrow -x_8 \implies$ no O3-plane, 1 O7-plane in D_8

• $\chi_{eff} = \chi(X) + 2 \int_X [O7] \wedge [O7] \wedge [O7]$ [Minasian, Pugh, Savelli]

Brane setup:

- Focus on brane setup which gives rise to **winding corrections**
- Stacks of branes around D_4 , D_6 and $D_2 = D_7 - D_1$
- $D7$ tadpole cancellation: $8[O7] \equiv 8([D_8]) = 16([D_2] + [D_4] + [D_6])$
- $D3$ tadpole cancellation:

$$N_{D_3} + \frac{N_{flux}}{2} + N_{gauge} = \frac{N_{O_3}}{4} + \frac{\chi(O7)}{12} + \sum_a \frac{N_a(\chi(D_a) + \chi(D'_a))}{48} = 38$$

→ Room for gauge and background three-form fluxes

- $D3$ -brane instanton on D_1
- Fluxes: $F_i = \sum_{j=1}^{h_{1,1}} f_{i,j} \hat{D}_j - \frac{1}{2} c_1(D_i) - \iota_{D_j}^* B \quad f_{1j} \in \mathbb{Z}$

$$B = \frac{1}{2} \hat{D}_1 + \frac{1}{2} \hat{D}_2 \quad \implies \quad \mathcal{F}_1 = \mathcal{F}_4 = \mathcal{F}_6 = 0; \quad \mathcal{F}_2 = \sum_{j=1}^{h_{1,1}} f_{2,j} \hat{D}_j \neq 0$$

Brane setup:

- $\mathcal{F}_2 \neq 0$

$$\left[\begin{array}{l} Sp(16) \rightarrow U(8) = SU(8) \times U(1) \\ q_{i2} = \int_X \hat{D}_i \wedge \hat{D}_2 \wedge \mathcal{F}_2; \quad q_{12} = -2f_{21}; \quad q_{42} = 2f_{26}; \quad q_{62} = 2f_{24}; \quad q_{72} = 0 \\ Re(f_2) = \frac{4\pi}{g_2^2} = \tau_2 - Re(S)h(\mathcal{F}_2) \quad h(\mathcal{F}_2) = \frac{1}{2}(f_{21}q_{12} + f_{24}q_{42} + f_{26}q_{62}) \\ \xi = \frac{1}{4\pi\mathcal{V}} \sum_{j=1}^{h_{1,1}} q_j 2t_j = \frac{1}{4\pi\mathcal{V}}(q_{12}t_1 + q_{42}t_4 + q_{62}t_6) \\ I_2^S = 2q_{12} - q_{42} - q_{62} \quad I_2^A = q_{42} + q_{62} \end{array} \right.$$

- No chiral intersection between $D7$ on D_2 and the instanton on D_1 : $I_{21} = q_{12} = 0$
- D-term constraint: $t_4 = \alpha t_6; \quad \alpha = -\frac{q_{62}}{q_{42}}$

Inflationary potential:

- D-term constraint: $t_4 = \alpha t_6; \quad \mathcal{V} = \frac{1}{\sqrt{2\alpha}} \sqrt{\tau_7} \tau_6 - \frac{1}{3} \tau_1^{3/2}$
- LVS stabilisation: $\langle \mathcal{V} \rangle = \frac{W_0}{4a_1 A_1} \sqrt{\langle \tau_1 \rangle} e^{a_1 \langle \tau_1 \rangle} \quad \langle \tau_1 \rangle = \left(\frac{3\hat{\xi}}{2} \right)^{2/3}$
- Winding 1-loop: $V_{g_s}^W = -2\kappa \frac{W_0^2}{\mathcal{V}^3} \sum_i \frac{C_i^W}{t_i^{(1)}} \quad \text{[Berg, Haack, Kors], [Berg, Haack, Pajer] [Cicoli, Conlon, Quevedo]}$
- HD $\alpha'{}^3 F^4$: $V_{F^4} = -\kappa^2 \frac{\lambda W_0^4}{g_s^{3/2} \mathcal{V}^4} \sum_{i=1}^{h_{1,1}} \Pi_i t_i; \quad \kappa = \frac{e^{K_{CS} g_s}}{8\pi} \quad \text{[Ciupke, Louis, Westphal]}$

$$V_{inf} = V_{g_s}^W + V_{F^4} = \kappa \frac{W_0^2}{\mathcal{V}^3} \left(\frac{A_1}{\tau_7} - \frac{A_2}{\sqrt{\tau_7}} + \frac{B_1 \sqrt{\tau_7}}{\mathcal{V}} + \frac{B_2 \tau_7}{\mathcal{V}} \right)$$

$$A_1 = \frac{3}{\pi} \frac{|\lambda| W_0^2}{\sqrt{g_s}}; \quad A_2 = C_W - \tilde{C}_W(\tau_7); \quad B_1 = \frac{(\alpha+1)}{\sqrt{2\alpha}} A_1; \quad B_2 = |C_3^W| - \hat{C}_W(\tau_7)$$

$$C_W = \sqrt{2\alpha} \left(C_1^W + \frac{C_2^W}{\alpha} \right); \quad \tilde{C}_W(\tau_7) = \frac{|C_4^W|}{(\alpha+1)} \sqrt{\frac{\alpha}{2}} \left(1 - \frac{\sqrt{2\alpha}}{(\alpha+1)} \sqrt{\frac{\tau_1}{\tau_7}} \right)^{-1}$$

$$\hat{C}_W(\tau_7) = \frac{C_5^W}{2} \left(1 + \frac{1}{\sqrt{2\alpha}} \frac{\tau_7^{3/2}}{\mathcal{V}} \right)^{-1} + \frac{C_6^W}{2} \left(1 + \sqrt{\frac{\alpha}{2}} \frac{\tau_7^{3/2}}{\mathcal{V}} \right)^{-1}$$

Single-field inflation - Requirements:

- $Vol_S^{1/4} \gg \sqrt{\alpha'}$ $\rightarrow \epsilon_{\tau_i} = \frac{1}{(2\pi)^4 g_s \tau_i} \ll 1$

- Dynamics inside KC:
$$\begin{cases} 2\alpha \langle \tau_1 \rangle < \tau_7 < \frac{V}{\sqrt{\langle \tau_1 \rangle}} & \alpha \geq 1 \\ \frac{2}{\alpha} \langle \tau_1 \rangle < \tau_7 < \frac{V}{\sqrt{\langle \tau_1 \rangle}} & \alpha \leq 1 \end{cases}$$

- Number of e-folds:
$$N_e \simeq 57 + \frac{1}{4} \ln(r_* V_*) - \frac{1}{3} \ln\left(\frac{V_{end}}{T_{rh}}\right)$$
$$T_{rh} \simeq 0.1 m_\phi \sqrt{\frac{m_\phi}{M_P}} \quad \text{[Cicoli, Piovano]}$$

- Amplitude of density perturbation matches COBE:
$$\frac{V_*^3}{V_*'^2} \simeq 2.6 \times 10^{-7}$$

- Check α' expansion:
$$\frac{\hat{\xi}}{2V} \ll 1$$

- EFT under control:
$$m_{inf} < H < m_{3/2} < M_{KK}^i < M_s < M_P; \quad M_{KK}^i = \frac{\sqrt{\pi} M_P}{\tau_i^{1/4} \sqrt{V}}$$

- Single field approximation:
$$\delta = \frac{H}{m_V} \sim \sqrt{\frac{V_*}{3V_{\alpha'}}} \lesssim 1$$

Single-field inflation - Parameters:

$$V = \kappa \frac{A_2 W_0^2}{\mathcal{V}^3 \sqrt{\langle \tau_7 \rangle}} \left(C_{dS} + c e^{-k\hat{\phi}} - e^{-\frac{k\hat{\phi}}{2}} + \mathcal{R}_1 e^{\frac{k\hat{\phi}}{2}} + \mathcal{R}_2 e^{k\hat{\phi}} \right)$$

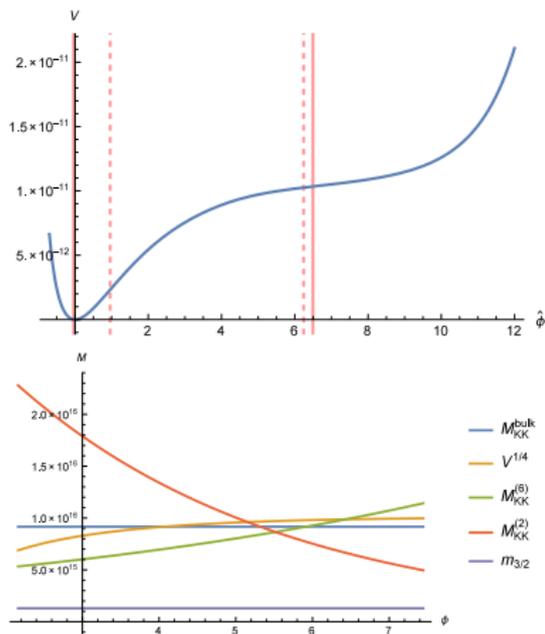
$$\mathcal{R}_1 = \frac{(\alpha+1)c}{\sqrt{2}\alpha} \frac{\langle \tau_7 \rangle^{\frac{3}{2}}}{\mathcal{V}} \ll 1; \quad \mathcal{R}_2 = \frac{(|C_3^W| - \hat{C}_W(\tau_7)) \langle \tau_7 \rangle^{\frac{3}{2}}}{(C_W - \hat{C}_W(\tau_7)) \mathcal{V}} \ll 1$$

$$c = \frac{3}{\pi(C_W - \hat{C}_W(\tau_7))} \frac{|\lambda| W_0^2}{g_s \langle \tau_7 \rangle} \sim \mathcal{O}(1);$$

[Grimm, Mayer, Weissenbacher]

Case 1: $ \lambda = 10^{-3}$, $\chi_{eff}(X)$	Case 2: $ \lambda = 10^{-7}$, $\chi(X)$
$\alpha = 1, \quad g_s = 0.114, \quad \mathcal{V} = 10^4, \quad W_0 = 80,$ $\langle \tau_1 \rangle = 1.91, \quad \xi = 0.067$ $C_1^W = C_2^W = 15, \quad C_3^W = 0.013,$ $ C_4^W = 18, \quad C_5^W = C_6^W = -5$	$\alpha = 1, \quad g_s = 0.25, \quad \mathcal{V} = 4500, \quad W_0 = 150,$ $\langle \tau_1 \rangle = 3.10, \quad \xi = 0.456$ $C_1^W = C_2^W = 0.034, \quad C_3^W = 10^{-5},$ $ C_4^W = 0.068, \quad C_5^W = C_6^W = -0.048$

Case 1: $|\lambda| = 10^{-3}$, $\chi_{eff}(X)$

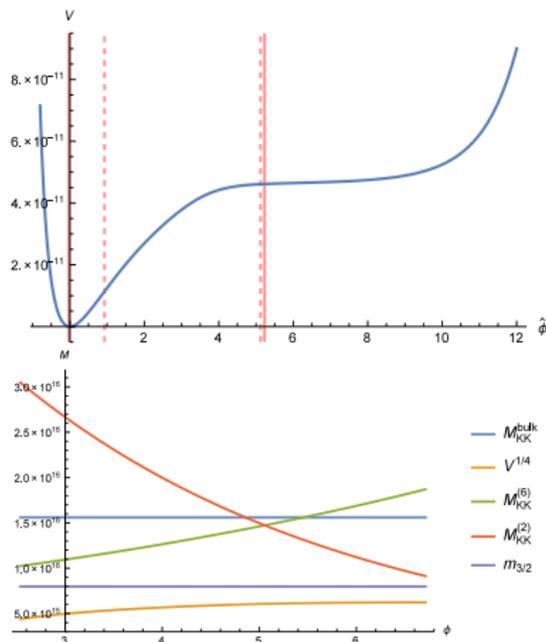


$$r \sim 10^{-3}, \quad n_s = 0.983, \quad \Delta N_{eff} = 0.39,$$

$$N_{max} \simeq 60, \quad m_\phi \simeq 4 \cdot 10^{13} \text{ GeV},$$

$$T_{rh} \simeq 1.8 \cdot 10^{10} \text{ GeV}, \quad N_e = 52, \quad \delta \simeq 1.6$$

Case 2: $|\lambda| = 10^{-7}$, $\chi(X)$



$$r \sim 0.0014, \quad n_s = 0.963, \quad \Delta N_{eff} = 0$$

$$N_{max} \simeq 58, \quad m_\phi \simeq 1.8 \cdot 10^{13} \text{ GeV},$$

$$T_{rh} \simeq 5.2 \cdot 10^9 \text{ GeV}, \quad N_e = 51, \quad \delta \simeq 0.05$$

Multi-field inflation - Equations of motion:

- Kinetic Lagrangian:

$$\mathcal{L}_{kin} = \frac{1}{2} \left(-\frac{\mathcal{V}'^2}{\mathcal{V}^2} + \frac{\mathcal{V}'\tau_7'}{\mathcal{V}\tau_7} - \frac{3\tau_7'^2}{4\tau_7^2} + \frac{\sqrt{\tau_1}\tau_7'\tau_1'}{2\mathcal{V}\tau_7} - \frac{\tau_1'^2}{4\mathcal{V}\sqrt{\tau_1}} \right)$$

- Potential:

$$V = \kappa \left[32A_s^2\pi^2 \frac{\sqrt{\tau_1}}{\mathcal{V}} e^{-4\pi\tau_1} - 8\pi A_s \frac{W_0\tau_1}{\mathcal{V}^2} e^{-2\pi\tau_1} + \frac{3\xi}{4g_s^{3/2}} \frac{W_0^2}{\mathcal{V}^3} + \frac{W_0^2}{\mathcal{V}^3} \left(\frac{A_1}{\tau_7} - \frac{A_2}{\sqrt{\tau_7}} + \frac{B_1\sqrt{\tau_7}}{\mathcal{V}} + \frac{B_2\tau_7}{\mathcal{V}} \right) + \frac{\delta_{up}}{\mathcal{V}^{4/3}} \right]$$

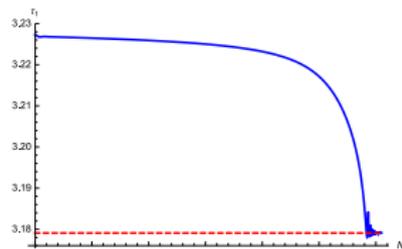
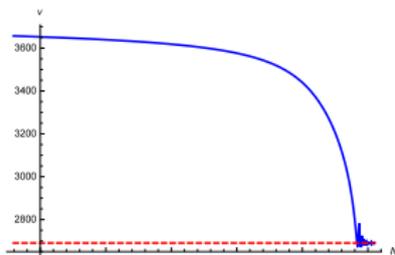
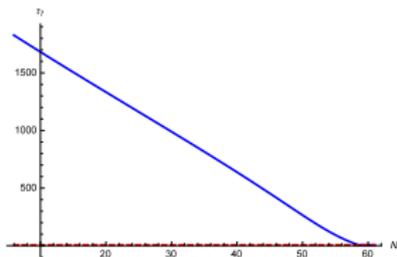
- EOM:

$$\begin{cases} \phi''^i + \Gamma_{jk}^i \phi'^j \phi'^k + (3 + \mathcal{L}_{kin}) \left[g^{ij} \frac{\partial_j V}{V} + \phi'^i \right] = 0 \\ H^2 = \left(\frac{\dot{a}}{a} \right)^2 = \frac{V}{3 + \mathcal{L}_{kin}} \end{cases}$$

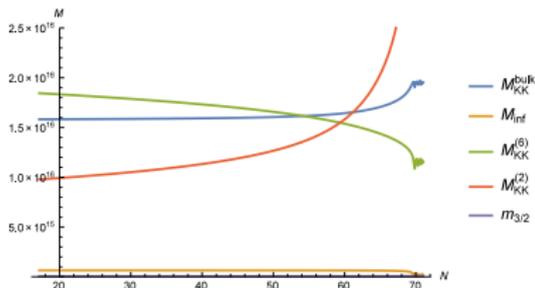
Multi-field inflation - Parameters and Results:

Case 1: $ \lambda = 10^{-3}$, $\chi(X)$	Case 2: $ \lambda = 10^{-6}$, $\chi(X)$
$\alpha = 1$, $g_s = 0.25$, $W_0 = 1$, $\xi = 0.456$, $A_s = 1 \cdot 10^4$, $C_1^W = C_2^W = 0.05$, $ C_3^W = 10^{-4}$, $ C_4^W = 0.1$, $C_5^W = C_6^W = -0.05$	$\alpha = 1$, $g_s = 0.25$, $W_0 = 50$, $\xi = 0.456$, $A_s = 6 \cdot 10^5$, $C_1^W = C_2^W = 0.05$, $ C_3^W = 10^{-4}$, $ C_4^W = 0.1$, $C_5^W = C_6^W = -0.05$
Min: $\mathcal{V} \simeq 3221$, $\tau_7 \simeq 6.4$, $\tau_1 = 3.18$	Min: $\mathcal{V} \simeq 2690$, $\tau_7 \simeq 6.5$, $\tau_1 = 3.18$
IC: $\mathcal{V} \simeq 4436$, $\tau_7 \simeq 2456$, $\tau_1 = 3.23$	IC: $\mathcal{V} \simeq 3670$, $\tau_7 \simeq 2036$, $\tau_1 = 3.22$
$n_s \simeq 0.967$, $r \simeq 0.002$, $N_{max} \simeq 58$ $\sqrt{P} \simeq 2 \cdot 10^{-7}$ vs $\sqrt{P_{\text{COBE}}} \simeq 2 \cdot 10^{-5}$	$n_s \simeq 0.97$, $r \simeq 0.002$, $N_{max} \simeq 69$ $\sqrt{P} \simeq 1 \cdot 10^{-5}$ vs $\sqrt{P_{\text{COBE}}} \simeq 2 \cdot 10^{-5}$

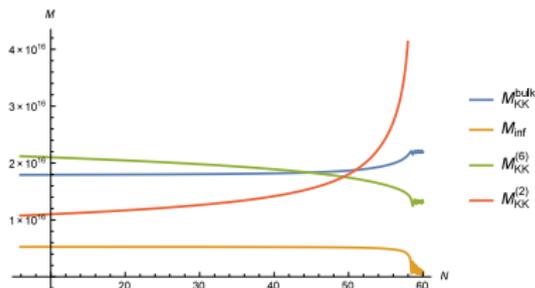
Multi-field inflation - Plots:



Case 1: $|\lambda| = 10^{-3}$, $\chi(X)$



Case 2: $|\lambda| = 10^{-6}$, $\chi(X)$



Conclusions:

- First explicit realisation of FI in concrete type IIB CY
- Inflationary parameters in good accordance with Planck results
- Observable tensor modes
- Kähler cone conditions strongly constrain the dynamics → hard to match COBE

Further investigations are needed:

- Compute χ_{eff} and λ in full detail
- Determination of the actual CY Kähler cone → **Shukla's Talk!**
- Find a mechanism to realise dS vacuum
- Generate density perturbations through curvaton mechanism? → **Pedro's Talk!**

Thank you for your attention