Global tensor matter transitions in F-theory

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F-theory [Vafa '96]

- Non-perturbative type IIB string theory
- Varying axio-dilaton
- Encoded in additional auxiliary torus over (compact complex) base manifold B via SL(2,Z) invariance
- Geometrization of physical properties
- Great framework for phenomenology



$$y^2 = x^3 + fx + g \rightarrow \operatorname{ord}(f, g, \Delta), \quad \Delta = 27g^2 + 4f^3$$

Elliptically-fibered Calabi-Yau n-folds (here n = 3)

F-theory (local and global)





- Codimension-one singularity induces non-Abelian gauge algebra
- Codimension-two singularity encodes matter
- No gravitational anomalies

- Mordell-Weil group dictates
 U(1) symmetries and non-Abelian gauge group
- Multi-sections lead to discrete gauge symmetries
- Gravitational anomalies

Non-flat fibers

 $\operatorname{ord}(f,g,\Delta) \ge (4,6,12)$

- Resolution in fiber not enough for smooth geometry
- Blow-up in the base
- Strongly coupled dynamics (superconformal matter)

Why do we care?

- Generically appear for GUT gauge groups, for SO(10) see e.g. [Buchmuller, MD, Oehlmann, Ruehle '17]
- Lead to non-trivial 6d SCFTs in decoupling limit, reviewed in [Heckman, Rudelius '18]



Strong coupling dynamics hard to study directly (no supergravity limit)

Indirect investigation via transitions

see also [Anderson, Gray, Raghuram, Taylor '15]



- 6d SUGRA defined by gauge algebra, matter spectrum and tensors
- Complex structure deformation to (4,6,12) point
- Deformation to tensor branch

6d anomalies

- Strict anomaly constraints in six dimensions
- Irreducible anomalies have to vanish

$$H - V + 29T - 273 = 0$$

$$\operatorname{tr}_{\mathbf{R}}F^{4} = B_{\mathbf{R}}\operatorname{tr}F^{4} + C_{\mathbf{R}}(\operatorname{tr}F^{2})^{2}$$

$$\operatorname{tr}_{\mathbf{R}}F^{2} = A_{\mathbf{R}}\operatorname{tr}F^{2}$$

$$\operatorname{tr}_{\mathbf{R}}F^{2} = A_{\mathbf{R}}\operatorname{tr}F^{2}$$

• Reducible can be cancelled by generalized Green-Schwarz mechanism, [Green, Schwarz '84], [Sagnotti '92], [Sadov '96]

$$\mathcal{I}_8^{\mathrm{red}} = -\frac{1}{32} \Omega_{\alpha\beta} X^\alpha X^\beta$$

Green-Schwarz and geometry

e.g. [Park '11], [Park, Taylor '11], [Morrison, Park '12]

$$\mathcal{I}_8^{\text{red}} = -\frac{1}{32}\Omega_{\alpha\beta}X^{\alpha}X^{\beta}, \quad X^{\alpha} = \frac{1}{2}a^{\alpha}\text{tr}R^2 - \frac{2}{\lambda}b^{\alpha}\text{tr}F^2 - \sum_{i,j}2b^{\alpha}_{ij}F_iF_j$$

- $\{H_{\alpha}\}$ a **basis** for $H_2(B,\mathbb{Z})$ then: $\Omega_{\alpha\beta} = H_{\alpha} \cdot H_{\beta}$ describes the participating tensor fields
- a^{α} is the vector describing the **anticanonical class** K_B^{-1}
- b^α is the vector describing the locus on which gauge algebra g is localized
- b_{ij}^{α} describes the coefficient of the height pairing $-\pi(\sigma_i \cdot \sigma_j)$ accounting for Abelian symmetries



- Divisor \mathcal{Z} with gauge algebra g
- Intersection with divisor D (possible matter locus)
- Complex structure **deformation to (4,6,12)** point
- Blow-up in base with self-intersection -1 curve E

Anomalies after blow-up

Additional tensor multiplet:

- Get rid of 29 degrees of freedom in hypermultiplet sector
- More possibilities for Green-Schwarz mechanism

Involved matter representations constrained?

- Non-Abelian representations fixed
- Abelian charges can be constrained

We use:

- Smoothness of fiber over E (no new gauge group)
- Properties of the **blow-down map** β :

 $\beta_*: H_2(\tilde{B}, \mathbb{Z}) \to H_2(B, \mathbb{Z}), \quad \beta_*^{-1}(D) \cdot \beta_*^{-1}(D')|_{\tilde{B}} = D \cdot D'|_B$

• Fixed transformation of anticanonical class:

$$K_{\tilde{B}}^{-1} = \beta_*^{-1}(K_B^{-1}) - E$$

- Orthogonality properties of Shioda map
- New basis of second homology:

$$\beta_*^{-1}(H_\alpha) = \tilde{H}_\alpha + h_\alpha E ,$$
$$\{H_\alpha\}_B \to \{\tilde{H}_\alpha, E\}_{\tilde{B}}$$

Anomalies in transition

From connection between geometry and anomalies:

$$\Omega_{\alpha\beta} \to \tilde{\Omega}_{AB} = \begin{pmatrix} \Omega_{\alpha\beta} - h_{\alpha}h_{\beta} & h_{\alpha} \\ h_{\beta} & -1 \end{pmatrix}$$
$$a^{\alpha} \to \tilde{a}^{A} = \begin{pmatrix} a^{\alpha} \\ a^{\alpha}h_{\alpha} - 1 \end{pmatrix}$$
$$b^{\alpha} \to \tilde{b}^{A} = \begin{pmatrix} b^{\alpha} \\ 0 \end{pmatrix}$$

Non-Abelian matter fixed

Back to matter states:

$$\begin{split} &\sum_{\mathbf{R}} \Delta n[\mathbf{R}] \operatorname{dim}(\mathbf{R}) < 29 \,, \\ &\sum_{\mathbf{R}} \Delta n[\mathbf{R}] B_{\mathbf{R}} = 0 \,, \\ &\frac{\lambda^2}{3} \sum_{\mathbf{R}} \Delta n[\mathbf{R}] C_{\mathbf{R}} = -1 \,, \\ &\frac{\lambda}{6} \sum_{\mathbf{R}} \Delta n[\mathbf{R}] A_{\mathbf{R}} = -1 \end{split}$$

- SU(5): $\Delta S \supset -(\mathbf{10} \oplus \mathbf{5} \oplus \mathbf{5} \oplus \mathbf{5})$
- SO(10): $\Delta S \supset -(\mathbf{16} \oplus \mathbf{10})$

Abelian matter

With assumption:

$$\Omega_{\alpha\beta}\Delta b^{\alpha}_{ij}(b^{\beta}-a^{\beta})=0$$

Curiously always satisfied for toric hypersurfaces with genus-one $\mathcal Z$

$$\lambda \sum_{\mathbf{R},i,j} \Delta n[\mathbf{R}, q_i, q_j] q_i q_j A_{\mathbf{R}} - \frac{1}{6} \sum_{i,j} \Delta n[q_i, q_j] q_i q_j = 0$$

Satisfied in all examples:

- SU(5): $\Delta S = -(\mathbf{10}_{-3/5} \oplus 3 \times \mathbf{5}_{1/5} \oplus 3 \times \mathbf{1}_1 \oplus \mathbf{1}_0)$
- SO(10): $\Delta S = -(\mathbf{16}_{1/4} \oplus \mathbf{10}_{1/2} \oplus \mathbf{1}_1 \oplus 2 \times \mathbf{1}_0)$

Toric geometry

e.g. [Klevers, Mayorga Pena, Oehlmann, Piragua, Reuter '14]

Using the **top construction** and connected **toric geometry:**

[Batyrev '94], [Candelas, Font '98], [Bouchard, Skarke '03]

Point in face corresponds to non-flat fiber see also [Braun, Grimm, Keitel '13]

$$h^{1,1}(Y_3) = \operatorname{rank}(G) + h^{1,1}(B) - 1 + \sum \delta_i n_{\mathrm{SCP}_i}$$

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Modification of number of neutral singlets in presence of nonflat fibers

- Consistency conditions for coupling strongly-coupled sectors to gravity see also [del Zotto, Heckman, Morrison, Park '15]
- Inclusion of global data controlling Abelian symmetries
- Toric realization in terms of hypersurfaces and tops

Outlook

- Extension to higher-order singularities
- Direct investigation of non-flat fibers [Achmed-Zade, Garcia-

Etxebarria, Mayrhofer '18]

 Comparison with Abelian flavor symmetries in decoupling limit [Lee, Regalado, Weigand '18]