
DAI-FREED ANOMALIES IN PARTICLE PHYSICS

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(Work in collaboration with Iñaki García-Etxebarria)

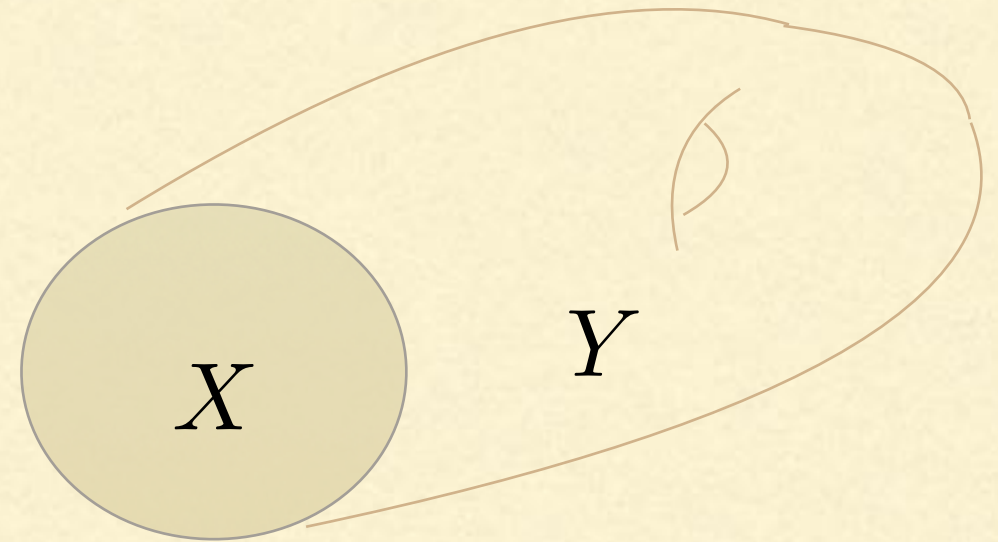
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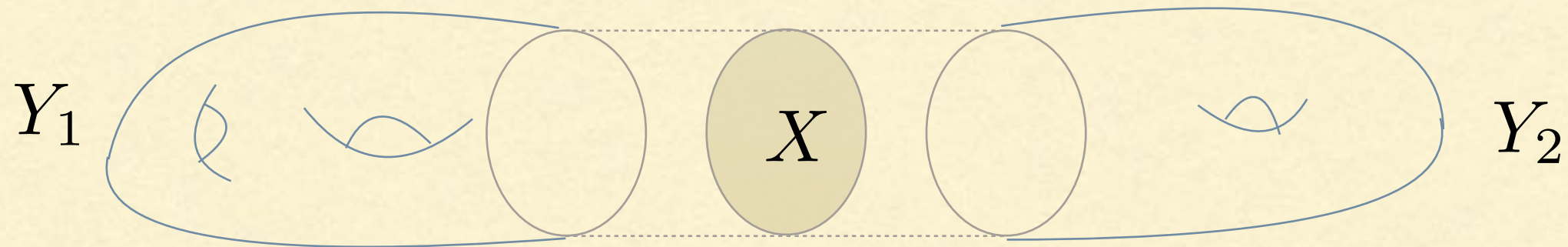
Utrecht University

QUICK RECAP

- Iñaki just explained the **Dai-Freed** formalism to compute anomalies of fermion systems in d dimensions.



- Anomalies cancel if $\exp(2\pi i \eta_Y)$ is independent of the choice of Y



- To ensure this, we must have **$\exp(2\pi i \eta_Y) = 1$** on any allowed $(d+1)$ manifold.

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This talk: Apply **Dai-Freed** to symmetries of interest in **particle physics**.

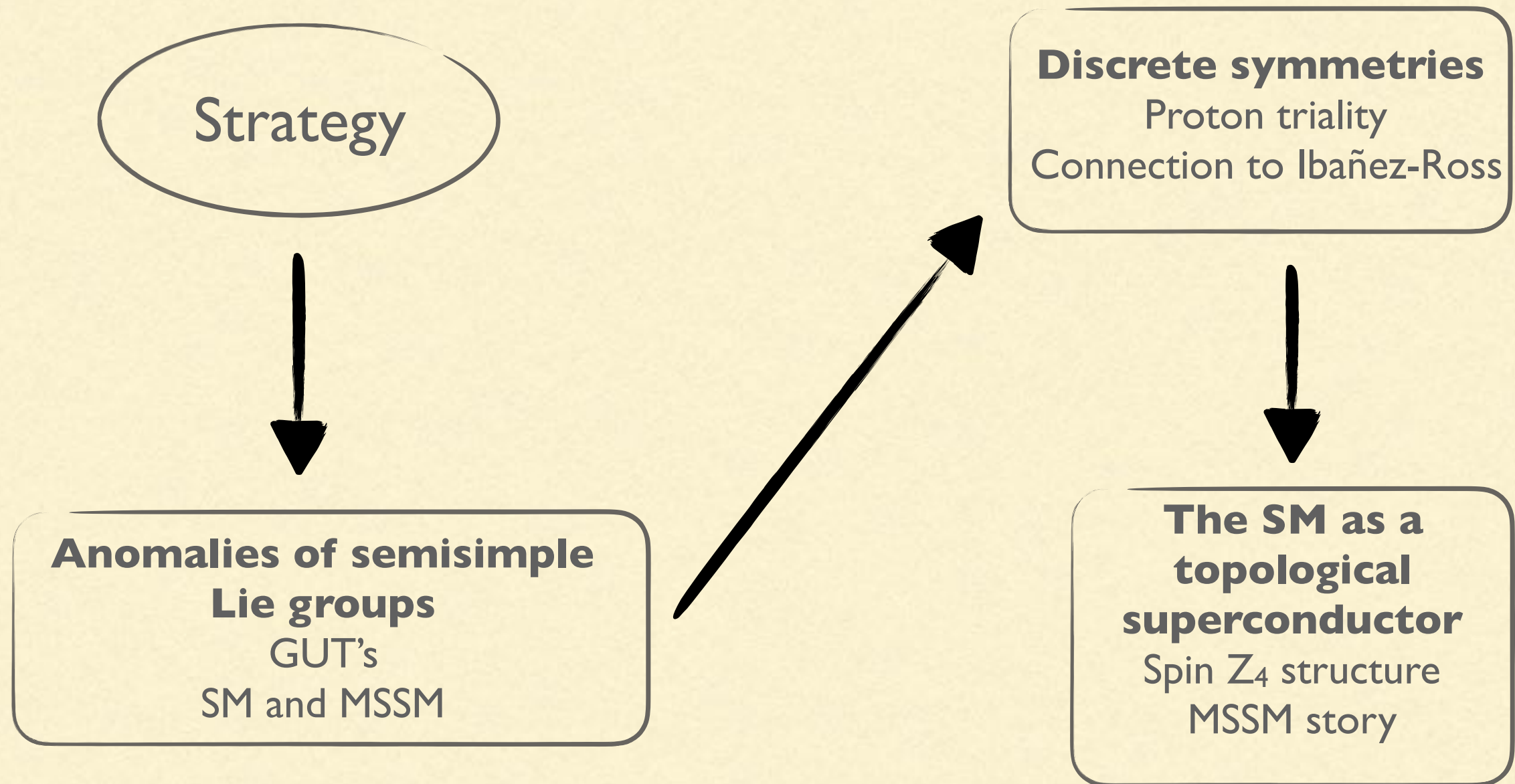
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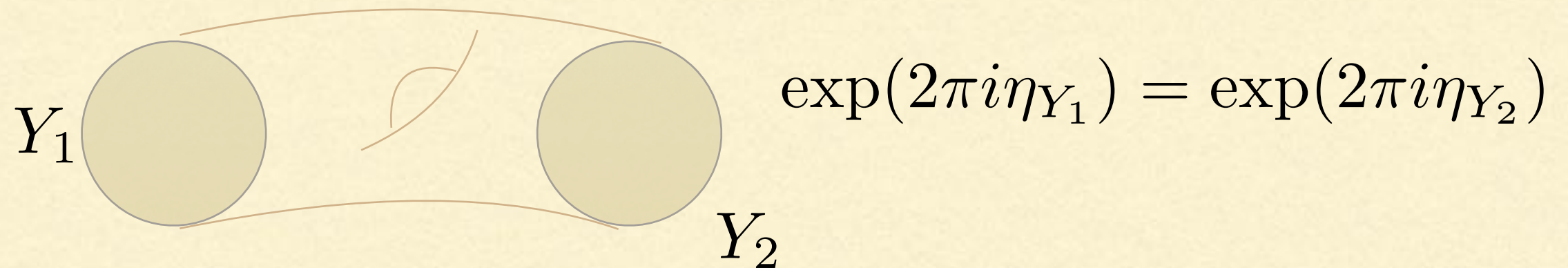
This talk: Apply **Dai-Freed** to symmetries of interest in **particle physics**.

Is the Standard Model Dai-Freed anomalous?

PLAN OF THE TALK



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- Once local anomalies cancel, η is a **bordism invariant**:



- **Bordism** is an equivalence relation, which defines **bordism groups**

$$\Omega_{d+1}^{Spin}(BG)$$

These classify $(d+1)$ -dimensional manifolds, with a principal G -bundle, modulo bordism (bundle extends over bordism too)

- Computed using **AHSS**.
 - η is a **group homomorphism** from the relevant bord. group to $U(1)$.
-

GENERAL STRATEGY

- Compute relevant bordism group
 - If it vanishes, there is no new anomaly.
 - Find a nontrivial manifold Y , compute η .
 - If it vanishes, there is no new anomaly.
 - If η is nonvanishing, there is an anomaly.
-

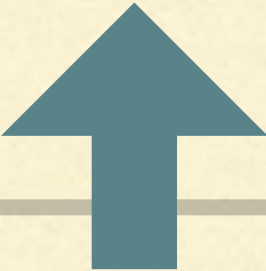
SEMISIMPLE LIE GROUPS

| G | $\Omega_{\mathrm{d}}^{\mathrm{Spin}}(\mathbf{BG})$ | | | | | | | | |
|------------------|----------------------------------------------------|----------------|----------------------------------|---|-------------------------------------------------|----------------|----------------|---|---------------|
| | 0 | 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 |
| $SU(2)$ | \mathbb{Z} | \mathbb{Z}_2 | \mathbb{Z}_2 | 0 | $2\mathbb{Z}$ | \mathbb{Z}_2 | \mathbb{Z}_2 | 0 | $4\mathbb{Z}$ |
| $SU(n > 2)$ | \mathbb{Z} | \mathbb{Z}_2 | \mathbb{Z}_2 | 0 | $2\mathbb{Z}$ | 0 | — | — | — |
| $USp(2k > 2)$ | \mathbb{Z} | \mathbb{Z}_2 | \mathbb{Z}_2 | 0 | $2\mathbb{Z}$ | \mathbb{Z}_2 | \mathbb{Z}_2 | 0 | $5\mathbb{Z}$ |
| $U(1)$ | \mathbb{Z} | \mathbb{Z}_2 | $\mathbb{Z}_2 \oplus \mathbb{Z}$ | 0 | $\mathbb{Z} \oplus \mathbb{Z}$ | 0 | — | — | — |
| $Spin(n \geq 8)$ | \mathbb{Z} | \mathbb{Z}_2 | \mathbb{Z}_2 | 0 | $2\mathbb{Z}$ | 0 | — | — | — |
| $SO(n \geq 3)$ | \mathbb{Z} | \mathbb{Z}_2 | $e(\mathbb{Z}_2, \mathbb{Z}_2)$ | 0 | $e(\mathbb{Z}, \mathbb{Z} \oplus \mathbb{Z}_2)$ | 0 | — | — | — |
| E_6, E_7, E_8 | \mathbb{Z} | \mathbb{Z}_2 | \mathbb{Z}_2 | 0 | $2\mathbb{Z}$ | 0 | 0 | 0 | $2\mathbb{Z}$ |
| G_2 | \mathbb{Z} | \mathbb{Z}_2 | \mathbb{Z}_2 | 0 | $2\mathbb{Z}$ | 0 | — | — | — |
| F_4 | \mathbb{Z} | \mathbb{Z}_2 | \mathbb{Z}_2 | 0 | $2\mathbb{Z}$ | 0 | 0 | 0 | — |

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THE STANDARD MODEL

- Experiments only probe the gauge **algebra** of the SM. There are four possibilities [Tong '17...]

$$\frac{SU(3) \times SU(2) \times U(1)}{\Gamma}, \quad \Gamma \in \{1, \mathbb{Z}_2, \mathbb{Z}_3, \mathbb{Z}_6\}$$

- The discrete group Γ acts trivially on the SM fermions.
 - $\Gamma = \mathbb{Z}_6$ is “maximal”: Includes bundles for any other choice of Γ .
 - This is also the group that embeds in $SU(5)$.
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- The SM fermion spectrum falls into **SU(5) representations**.
 - Any **(SU(3) x SU(2) x U(1))/Z₆** bundle is a **SU(5) bundle** too!
 - As far as anomalies are concerned, the SM is **equivalent** to the SU(5) GUT. But since

$$\Omega_5^{Spin}(BSU(5)) = 0$$

The SM is free of Dai-Freed anomalies

- Similar situation for Spin(10).
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-
- To get an anomaly, we need to look at **more general** spaces.
 - What about the SM in **non-orientable spaces**?
 - Only makes sense if one assumes CP breaking in SM is **spontaneous**.
 - Need a Pin structure to define fermions, which can change cob. groups, e.g. $\Omega_6^{\text{Pin}^-} = \mathbb{Z}_{16}$, but $\Omega_6^{\text{Spin}} = 0$.
 - Majorana masses require a Pin⁺ structure [Berg et al '00].
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 - Majorana masses require a Pin⁺ structure [Berg et al '00].
 - We have again

$$\Omega_5^{\text{Pin}^+}(BSU(5)) = \Omega_5^{\text{Pin}^-}(BSU(5)) = 0$$

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- Last try: SM + right-handed neutrinos + gauged (B-L).
 - Since all fermion charges under (B-L) are odd, we can now consider the SM on **Spin^c manifolds** or even **Pin^c manifolds**
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 - Since all fermion charges under (B-L) are odd, we can now consider the SM on **Spin^c manifolds** or even **Pin^c manifolds**
 - Still,

$$\Omega_5^{Spin^c}(BSU(5)) = \Omega_5^{Pin^c}(BSU(5)) = 0$$

so we find no anomalies in the SM.

DISCRETE CYCLIC GROUPS

- Lucky! Bordism groups & η invariants already computed by mathematicians [Bahri-Gilkey '87, Gilkey '89, Gilkey-Botvinnik '94] both for Spin and Spin^c cases. They are **nontrivial**.
- We can compare with **known anomalies** of discrete symmetries. [Ibañez-Ross' 91]. These were originally obtained by demanding that the Z_n embeds in a U(1).

$$2 \sum s_i \equiv 0 \bmod N, \quad \sum s_i^3 \equiv 0 \bmod N$$

- Only **linear** constraints are UV-independent [Banks-Dine' 91]
-

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- There is a nontrivial **Dai-Freed** anomaly coming from evaluating the η in a generalized lens space [Bahri-Gilkey '86, Gilkey '89, Gilkey-Botvinnik '94].
 - We get constraints which are cubic in the charges: A “remnant” of the **cubic Ibáñez-Ross constraint**

$$\sum -4s_i^3 + (N^2 + 3)s_i \equiv \text{mod } 24N$$

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$$\sum -4s_i^3 + (N^2 + 3)s_i \equiv \text{mod } 24N$$

- Are these UV-sensitive? **YES**
 - Topological GS term that can forbid some of the bundles [Ibáñez '92, Garcia-Etxebarria-Hayashi-Ohmori-Tachikawa-Yonekura '17]
-

- For Z_3 ,

$$\sum s_i \equiv 0 \pmod{3} \quad (\text{Linear IR})$$

$$\sum s_i \equiv 0 \pmod{9} \quad (\text{Dai-Freed})$$

This one has phenomenological consequences: Proton triality (and also hexality) in the MSSM is a IR-anomaly free Z_3 symmetry, but it has a mod 9 anomaly

| Q | u | d | l | e | H | \bar{H} |
|-----|-----|-----|-----|-----|-----|-----------|
| 0 | -1 | 1 | -1 | -1 | 1 | -1 |
| 0 | -2 | -5 | -5 | 1 | 5 | 5 |

$$\sum_{MSSM} s_i = 3 \pmod{9}$$

- Dai-Freed anomaly cancellation requires 3k generations.
- Consistent with previous results [Dreiner et al. '04]: $U(1)$ embedding of proton triality only with gen. dependent charges.

SM = TOP. SUPERCONDUCTOR

- **Topological superconductor:** 1st example of Dai-Freed anomaly [Kapustin-Thorgren-Turzillo-Wang '14, Witten 015, Hsieh-Cho-Ryu '16]

- T-invariant 3d fermions. Global grav. anomaly requires multiple of 8.

- Dai-Freed enhances to a multiple of 16, because

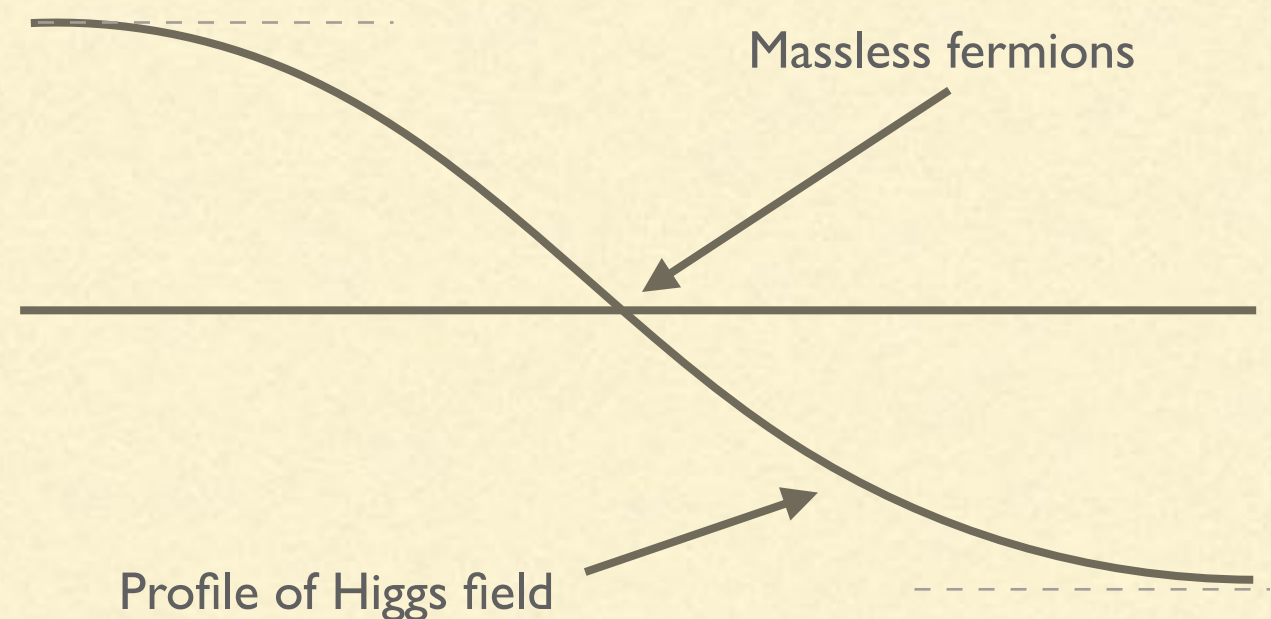
$$\Omega_4^{Spin^+} = \mathbb{Z}_{16}$$

- # of fermions/ generation in SM + rh neutrinos = 16. **Not a coincidence!**
-

- The SM + rh neutrinos has a Z_4 symmetry (center of $\text{Spin}(10)$) that acts on every fermion by multiplication by i .
- We can use this to put the SM on manifolds with a $\text{Spin}^{\mathbb{Z}_4}$ structure [Tachikawa-Yonekura '18]. Transition functions of the spinors in $(\text{Spin} \times Z_4)/Z_2$

- The **Smith homomorphism** maps

$$\Omega_{d+1}^{\text{Spin}^{\mathbb{Z}_4}} \rightarrow \Omega_d^{\text{Pin}^+}$$



- Physical interpretation: Higgsing the Z_4 w. a nontrivial bundle, there is a 3d locus with massless Pin^+ fermions.

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 - # of fermionic superpartners:
 - 12 gauginos (8+3+1)
 - 4 higgsinos
 - Total: 16! **Dai-Freed anomaly vanishes.**
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 - # of fermionic superpartners:
 - 12 gauginos ($8+3+1$)
 - 4 higgsinos
 - Total: 16! **Dai-Freed anomaly vanishes.**
 - Only works because of the detailed structure of SM: Dim. of gauge group + EWSB sector.
 - No obvious relation to GUT's. Related to reflections of compactification manifold? [Tachikawa-Yonekura '18]
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CONCLUSIONS

- We've explored a new kind of anomaly in four dimensional gauge theories of phenomenological interest.
 - SM and GUT's are anomaly free. Can put SM on non-Spin manifolds (related to topological superconductor).
 - New anomalies for discrete symmetries e.g. proton triality.
 - Outlook
 - We only checked a few theories!
 - Is the Z_4 in the MSSM telling us something?
-

DZIĘKI!
