# Group Theory Aspects of Non-Supersymmetric Heterotic Partition Functions 

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## Introduction

- SUSY breaking string models have received a lot of interest in the past decades
- In particular, there is the $\mathrm{SO}(16) \times \mathrm{SO}(16)$ heterotic string with $\mathcal{N}=0$ in 10 dimensions
- Alternatively, one may break SUSY during compactification: e.g. Scherk-Schwarz, or particular orbifold geometries


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In general, heterotic non-SUSY string vacua suffer from a bunch of problems, such as a too large cosmological constant and instabilities

## Cosmological constant/dilaton tadpole

As one can see from general arguments, the non-SUSY heterotic partition function $\mathcal{Z}(\tau, \bar{\tau})$ is nonzero

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\Lambda \sim \int \frac{\mathrm{d}^{2} \tau}{\tau_{2}^{2}} \mathcal{Z}(\tau, \bar{\tau})
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There are various proposals for this to vanish

- (space-time) Supersymmetry
- (generalized) Atkin-Lehner symmetry


## Heterotic Orbifolds

Orbifolds can be obtained in two subsequent steps

- define and mod out a lattice $\left\{e_{\alpha}\right\}_{\alpha=1, \ldots, D}$ to obtain a torus
- mod out a discrete isomorphism of the lattice (point group $P$ )

For Abelian point groups
$\mathrm{SO}(6) \supset D_{6}(g)=\left(\begin{array}{ccc}\mathrm{e}^{2 \pi \mathrm{i} v_{1} J_{12}} & & \\ & \mathrm{e}^{2 \pi i \mathrm{v}_{2} /_{34}} & \\ & & \mathrm{e}^{2 \pi \mathrm{i} v_{3} J_{56}}\end{array}\right) \rightarrow v_{g}=\left(v_{1}, v_{2}, v_{3}\right)$
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$\mathcal{N}=1$ SUSY $\leftrightarrow \sum v_{i}=0 \bmod 2$

Heterotic Orbifolds arise by an embedding of the geometric rotations and translations into the $\mathrm{E}_{8} \times \mathrm{E}_{8}$ d.o.f.

## The Heterotic Partition Function

The partition function can be organized

$$
\mathcal{Z}=\sum_{g, h} \mathcal{Z}\left[\begin{array}{l}
g \\
h
\end{array}\right] \quad \text { where } g, h \in \text { space group }
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It is known (via generalizations of the Jacobi Abstruse Identity):
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However, the unbroken SUSY transformations need not be the same for all $(g, h)$ pairs

Local "SUSY enhancement": An $\mathcal{N}=2$ Example
$\mathrm{A} \mathbb{Z}_{2} \times \mathbb{Z}_{2}$ orbifold with roto-translation
Fischer, Ratz, Torrado, Vaudrevange ' 12

$$
\begin{aligned}
v_{\theta} & =\left(0, \frac{1}{2},-\frac{1}{2}\right), \quad v_{\omega}=\left(\frac{1}{2},-\frac{1}{2}, 0\right) \\
\text { with generators } \quad g_{\theta} & =\left(\theta \left\lvert\, \frac{1}{2} e_{2}\right.\right), \quad g_{\omega}=(\omega \mid 0)
\end{aligned}
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with generators $\quad g_{\theta}=\left(\theta \left\lvert\, \frac{1}{2} e_{2}\right.\right), \quad g_{\omega}=(\omega \mid 0)$
Because of the roto-translation, $\left[g_{\theta}, g_{\omega}\right] \neq 0$, and hence

$$
\mathcal{Z} \not \supset \mathcal{Z}\left[\begin{array}{l}
g_{\theta} \\
g_{\omega}
\end{array}\right] \quad \text { (the same holds for all } \mathrm{SL}(2, \mathbb{Z}) \text { images) }
$$

## $\mathcal{N} \geq 2$ SUSY in each ( $g, h$ )-twisted sector,

 but$\mathcal{N}=1$ in the intersection!

Q: Can one achieve the same for $\mathcal{N}=0 \rightarrow \mathcal{N} \geq 1$ ?

## Four dimensional spinor representation

On the space-time fermions, each $D_{6}(g) \in 6$ of $\mathrm{SO}(6)$ corresponds to a $D_{4}(g) \in 4$

Because Spin $(6) \cong \operatorname{SU}(4)$

## 4 needs to be in $\operatorname{SU}(4)$, i.e. $\operatorname{det} D_{4}(g)=1$ for all $g$

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Condition 1

## Caveats

- Elements of the geometric space group have more than one embedding into spinor space $\rightarrow$ Witten twist
- Possible degeneracies lead to spinor embeddings that are non-isomorphic to the geometric point group

For our considerations: do not care too much about whether spinor embedding corresponds to geometric embedding (just look at all possible actions on spinor space)

## Realizing $\mathcal{N}=0$

Every element of the point group commutes with the identity, therefore the untwisted partition function is given by

$$
\mathcal{Z}_{\text {untw. }}=\sum_{g \in P} \mathcal{Z}\left[\begin{array}{l}
\mathbb{1} \\
g
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If we insist on SUSY breaking, then the projector of the untwisted sector must have trace 0

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or in terms of the representation content
4 does not contain a trivial singlet of $P$

## Interplay: Modular orbits vs. projectors

Let us for the moment forget about the space group and consider the point group only
Then, all constructing elements $g$ are just elements of a finite group, i.e. $\exists N$ s.t. $g^{N}=\mathbb{1}$

$$
\mathcal{Z}\left[\begin{array}{l}
\mathbb{1} \\
g
\end{array}\right] \xrightarrow{S} \mathcal{Z}\left[\begin{array}{l}
g \\
\mathbb{1}
\end{array}\right] \xrightarrow{T} \mathcal{Z}\left[\begin{array}{l}
g \\
g
\end{array}\right] \xrightarrow{T^{N-1}} \mathcal{Z}\left[\begin{array}{l}
g \\
\mathbb{1}
\end{array}\right]
$$

There may exist sectors that are not connected to these orbits, e.g. in Abelian $\mathbb{Z}_{N} \times \mathbb{Z}_{M}$ orbifolds.

We always assume these sectors can be removed

## Conditions for $\mathcal{N} \geq 1$ locally

We are effictively dealing with projectors of the type

$$
\sum_{n=0}^{N-1} \mathcal{Z}\left[\begin{array}{c}
g \\
g^{n}
\end{array}\right] \xrightarrow{\text { in group theory language }} \frac{1}{N} \sum_{n=0}^{N-1} \operatorname{tr} D_{4}\left(g^{n}\right)
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In representation theory language, this amounts to requiring that for all $g$

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Then our rationale becomes
$\mathcal{Z}\left[\begin{array}{l}\mathbb{1} \\ g\end{array}\right]$ and entire modular orbit vanishes


## A group-theoretical conjecture

Conjecture. For a given discrete group $G$, there does not exist a four dimensional representation with the properties that
(i) it has determinant 1 ,
(ii) it does not contain a trivial singlet of $G$,
(iii) to every $\mathbb{Z}_{N}$ subgroup of $G$ it branches into at least one trivial singlet of that $\mathbb{Z}_{N}$ subgroup.

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We have tested this conjecture for all point groups relevant for toroidal orbifolds, and for $\mathcal{O}(100000)$ more discrete groups.

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We have tested this conjecture for all point groups relevant for toroidal orbifolds, and for $\mathcal{O}(100000)$ more discrete groups.

It is crucial to insist on det $=1$ and dimension 4 !

## A chance for (generalizations of) Atkin-Lehner symmetry?

Idea: if $\mathcal{Z}(\tau, \bar{\tau}) \neq 0$, it might still integrate to zero

- Guaranteed if e.g. $\mathcal{Z}\left(-\frac{1}{N \tau},-\frac{1}{N \tau}\right)=-\mathcal{Z}(\tau, \bar{\tau})$ known as Atkin-Lehner symmetry Moore '87
- does not seem to work (at least not with four non-compact dimensions) TR Taylor '87, Balog and Tuite ' 87
but:
- We see that a large portion of the $\mathcal{Z}\left[\begin{array}{l}\mathbb{1} \\ g\end{array}\right]$ vanishes even if $\mathcal{N}=0$ globally
- Hence, not the entire partition function has to have a certain symmetry


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## Conclusions

- we have translated the vanishing condition for $(g, h)$-twisted sectors to a condition for the branching of the 4 to $\mathbb{Z}_{N}$ subgroups
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- the partial results point towards alternative solutions to the cosmological constant problem in (heterotic) non-SUSY strings


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