## Multi-field effects in a simple extension of ${\bf R}^2\, \text{Inflation}$

Taro Mori (KEK/SOKENDAI)

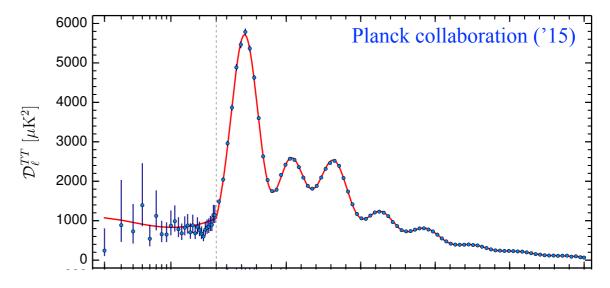
based on JCAP 1710 (2017) No.10, 044/arXiv:1705.05638 with Kaz Kohri(KEK) and Jonathan White

StringPheno18@Univ. of Warsaw 4/7/2018

### Introduction

### **\*Inflation**

- Horizon problem
- Flatness problem
- Origin of large scale structure

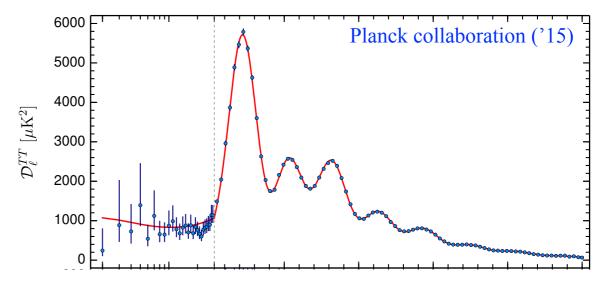


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### **\*Motivation of Multi-field**

- Supergravity or Superstring theory may provide many scalar fields
- → Multi-field models
- In multi-field models, curvature perturbation is NOT conserved on super-horizon scale.
- →Any feature which single-field models don't have?

### Set up: A simple extension of Starobinsky model

van de Bruck and Paduraru ('15)

### Jordan frame

$$S_{J} = \int d^{4}x \sqrt{-g} \left[ \frac{m_{pl}^{2}R}{2} + \frac{\mu}{2}R^{2} \right] + \int d^{4}x \sqrt{-g} \left[ -\frac{1}{2}g^{\mu\nu}\partial_{\mu}\chi\partial_{\nu}\chi - \frac{1}{2}m_{\chi}^{2}\chi^{2} \right]$$

### **Motivation:**

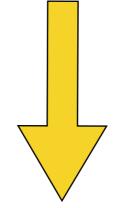
- Check the multi-field effects on asymptotic potential
- Essentially same with a class of models based on Supergravity Ellis et al ('13-16)

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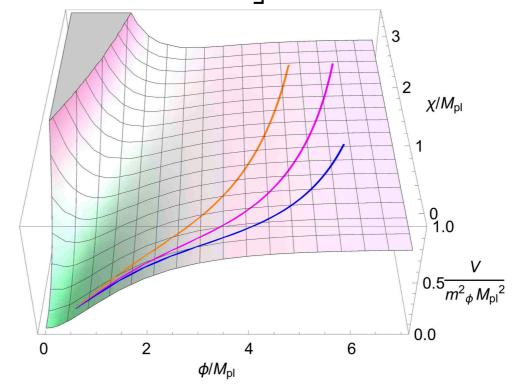
conformal transformation : where 
$$\tilde{g}_{\mu\nu}=\Omega^2 g_{\mu\nu} \ \ \Omega^2=1+\frac{2\mu R}{m_{pl}^2}$$

### **Einstein frame**

$$S_E = \int d^4x \sqrt{-\tilde{g}} \left[ \frac{m_{pl}^2}{2} \tilde{R} - \frac{1}{2} (\partial \phi)^2 - \frac{1}{2} e^{-2\alpha\phi} (\partial \chi)^2 - \tilde{V} \right]$$

$$e^{2\alpha\phi} = 1 + \frac{2\mu R}{m_{pl}^2} \quad \alpha = \frac{1}{\sqrt{6}m_{pl}}$$

$$\tilde{V} = \frac{m_{pl}^4}{8\mu} (1 - e^{-2\alpha\phi})^2 + \frac{1}{2} m_{\chi}^2 e^{-4\alpha\phi} \chi^2$$



### **Curvature perturbation**

### **\*Spectrums**

$$\langle \zeta_{\mathbf{k_1}} \zeta_{\mathbf{k_2}} \rangle = (2\pi)^3 \delta(\mathbf{k_1} + \mathbf{k_2}) P_s(k_1)$$
$$\langle \zeta_{\mathbf{k_1}} \zeta_{\mathbf{k_2}} \zeta_{\mathbf{k_3}} \rangle = (2\pi)^3 \delta(\mathbf{k_1} + \mathbf{k_2} + \mathbf{k_3}) B_s(k_1, k_2, k_3)$$

$$n_s \equiv \frac{d \log P_s}{d \log k} \qquad r \equiv \frac{P_t}{P_s}$$

$$f_{\rm NL} \equiv \frac{B_s(k_1, k_2, k_3)}{P_s(k_1)P_s(k_2) + \text{cyc. perm.}}$$

Lyth, Malik and Sasaki ('04)

### In single-field cases…

$$\dot{\zeta} \propto \left(\frac{k}{aH}\right)^2 \sim 0$$

### conserved on super-horizon scale

### **Consistency relation:**

in the squeezed limit

$$f_{\rm NL} \sim (1 - n_s) \sim \mathcal{O}(10^{-2})$$

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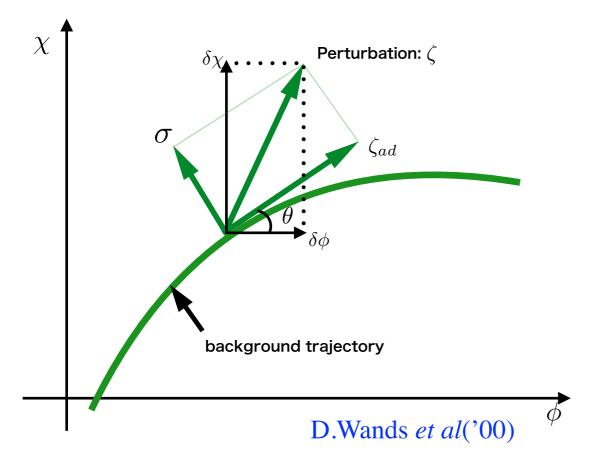
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### **\*Multi-field effects**

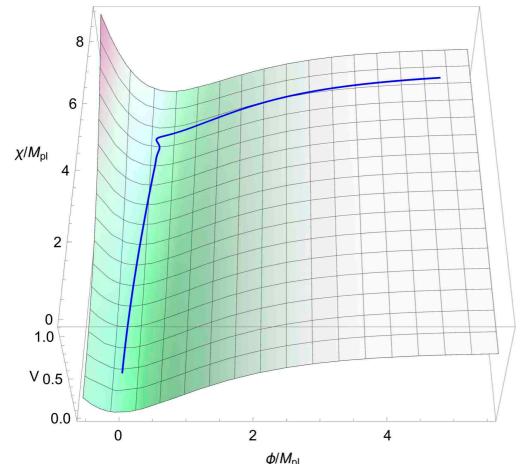


isocurvature mode can transfer to adiabatic mode, **¿** evolves with time on super-horizon scale

- →We need to take into account multi-field effects on each quantities.
- \*  $f_{NL}$  can be enhanced by multi-field effects?

### Superhorizon evolution of perturbations

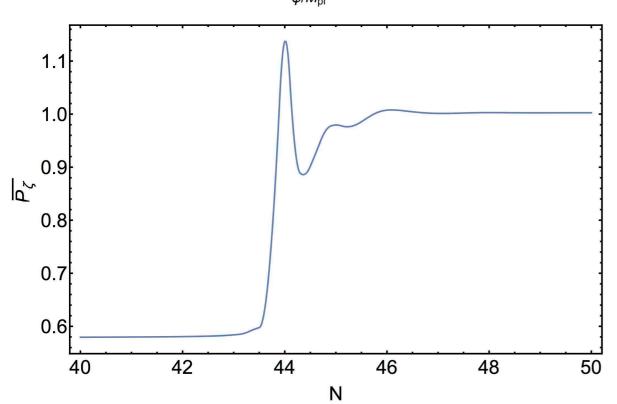
$$m_{\chi}/m_{\phi} = 0.1$$

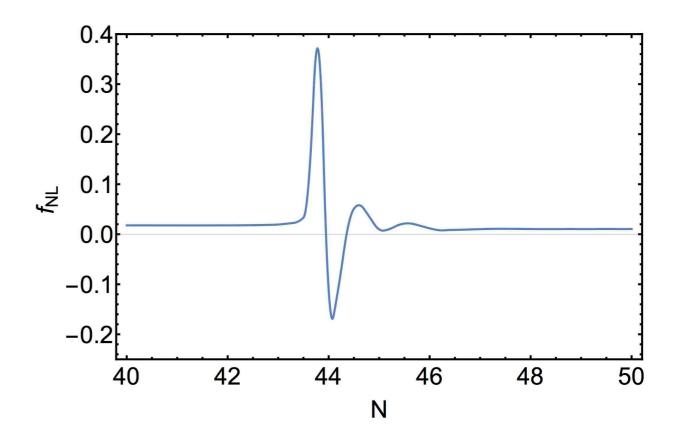


### We developed numerical approach based on $\delta N$ -formalism

If a trajectory turns/oscillates, each quantities may have steps, peaks or oscillations.

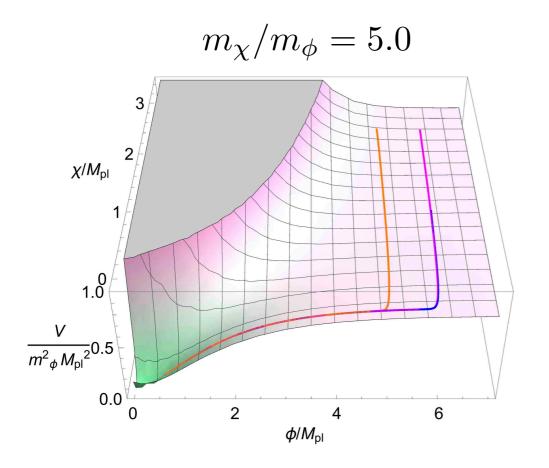
### **Evolutions on supuerhorizon scale!**

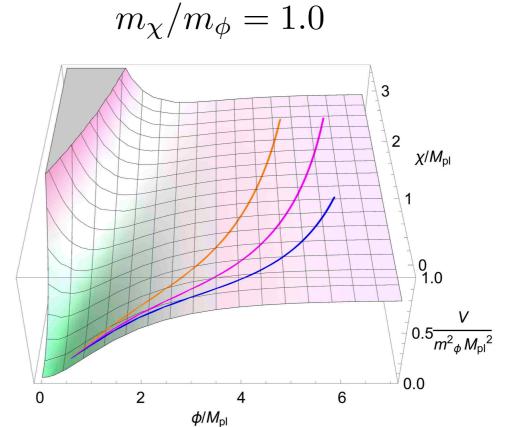


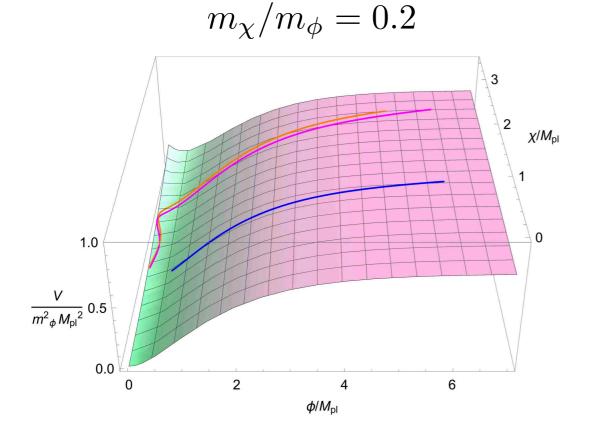


Field space search  $m_{\chi}/m_{\phi} = 1.0$  $n_s$  $n_s = 0.968 \pm 0.006$  $N_{
m total}$ (68% C.L.) 15 15 80 1.0 10 NX 10  $\chi M_{
m pl}$ 40 20 5 5 8.0 0.7 0 0, 2 3 5 4 6 2 3 4 5 6 1  $\phi/M_{\rm pl}$  $\phi/M_{\rm pl}$  $f_{NL} = 0.8 \pm 5.0$ r < 0.11 $f_{NL}$ 68% C.L. 95% C.L. 15 15 10 NX 10 0.003  $\chi M_{
m pl}$ 0.001 0.025 5 5 0.050 0.05 0.250 0.1 0, 0 2 3 2 3 4 5 6 5 4 6  $\phi/M_{\rm pl}$  $\phi/M_{\rm pl}$ 

### **Various mass-ratios**





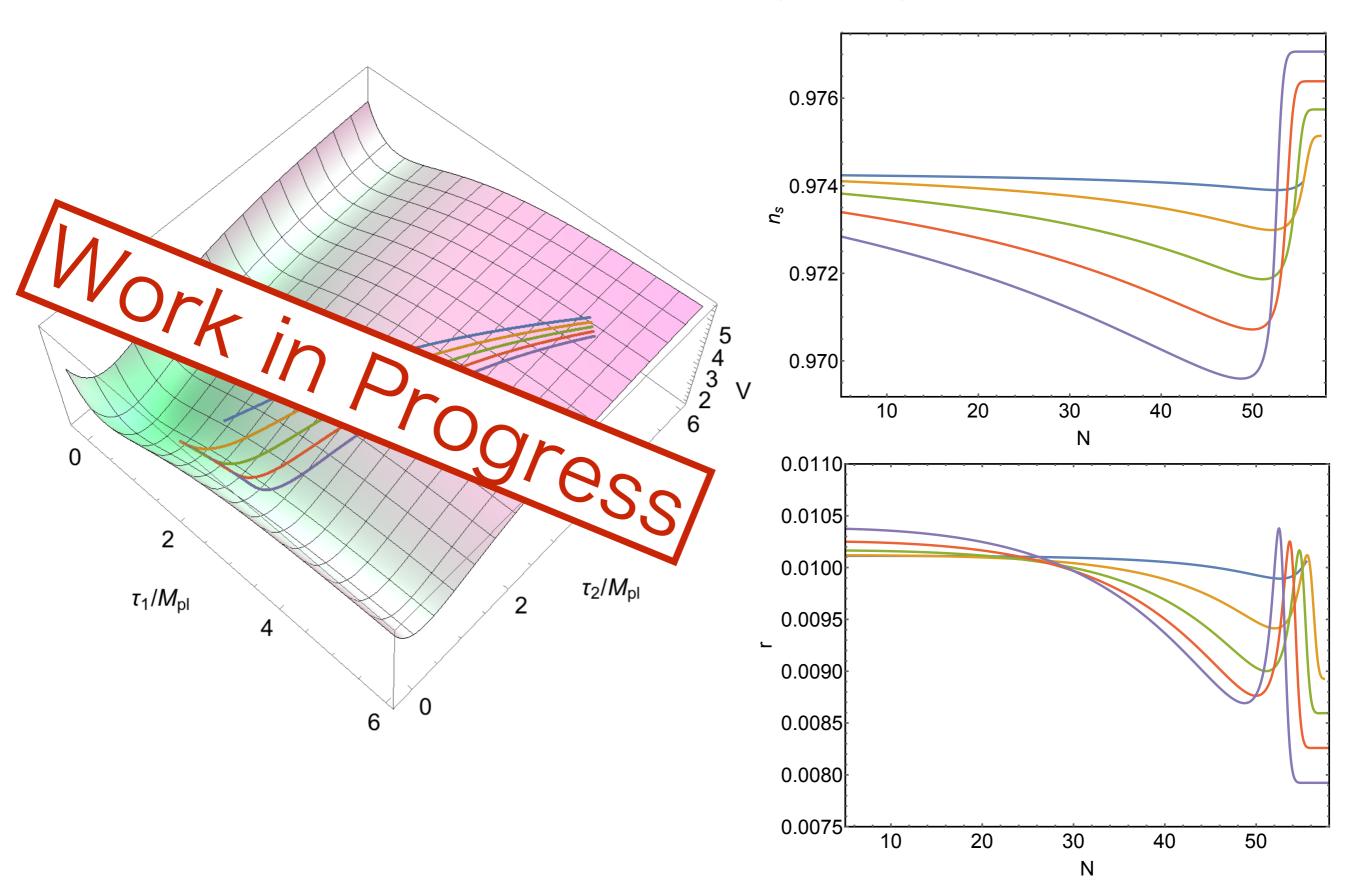


### We investigated the mass-ratios

$$10^{-3} \le m_{\chi}/m_{\phi} \le 10^3$$

### Multi-field model from string compactification

with Joe Conlon(Oxford), Kaz Kohri 18XX.XXXXX



### Summary

- There is a strong motivation to consider multi-field inflation models in string/supergravity inspired model-building.
- In the case of multi-field, we need to take into account transfer of isocurvature perturbation.
- In this work, we established the method to compute predictions in interacting multi-field models with non-canonical kinetic terms.
- $\cdot f_{NL}$  is as small as single field cases.

\*If you are interested in multi-field analysis,

I am happy to discuss with you!

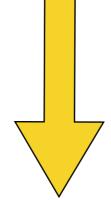
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### This model contains...

- Two inflatons
- Interaction term
- Non-canonical kinetic term

### $\delta N$ -formalism

Sasaki, Stewart('96) Wands et al ('00)

### On super-horizon scale,

$$\zeta \sim \delta N \equiv \mathcal{N}(t_*, t; \mathbf{x}) - N(t_*, t)$$
$$= N_I \delta \phi^I + \frac{1}{2} N_{IJ} \delta \phi^I \delta \phi^J$$

I : derivatives wrt  $\phi^I$ 

To compute each quantities, we only need background dynamics and  $N_{I},\ N_{I,J}$ 

### \*Non-canonical kinetic terms: Curved field space

 $\mathcal{L}_{kin} = -\frac{1}{2} \mathcal{G}_{IJ} \partial^{\mu} \phi^{I} \partial_{\mu} \phi^{J}$ 

Yokoyama et al('08)

Elliston et al('12), Kaiser et al('13)

 $\mathcal{G}_{IJ}$  can be interpreted as field space metric. In curved(non-trivial) field space, covariant formalism is useful.

### **Points:**

- 1, Contracting with field space metric
- 2, Derivatives should be replaced with Covariant derivatives

$$\partial_I A_J$$

$$\downarrow$$

$$\mathcal{D}_I A_J = \partial_I A_J - \Gamma_{IJ}^K A_K$$

### $\delta N$ -formalism

### **Power Spectrum:**

$$P_s = \frac{H_*^2}{4\pi^2} N^I N_I, \qquad P_t = \frac{H_*^2}{4\pi^2}$$

### spectral tilt:

$$n_s = 1 - 2\epsilon_* - 2\frac{1 + N_A(\frac{1}{3}R^{ABCD}\frac{V_BV_C}{V^2} - \frac{V^{;AD}}{V})N_D}{N^I N_I}$$

### tensor-scalar ratio:

$$r = \frac{P_t}{P_s} = \frac{1}{N^I N_I}$$

### **Non-Gaussianity:**

$$f_{NL} = \frac{5}{6} \frac{N^I N^J \mathcal{D}_I \mathcal{D}_J N}{(N^K N_K)^2}$$