Cosmic Archaeology with Gravitational Wave Signals from Cosmic Strings

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Based on:

Y. Cui, D. Morrissey, ML and J. D. Wells. arXiv:1711.03104, arXiv:1807.XXXXX





- A brief review of cosmic strings
- Stochastic GW background from Cosmic Strings
- Cosmic Archaeology
 - GW frequency temperature relation
 - Detection capabilities
 - Non standard cosmologies
- Conclusions

• What are cosmic strings?

 $\rightarrow~$ Stable one-dimensional topological defects

• The origins of cosmic strings:

- \rightarrow Prediction from Superstring theory: (F-) string, D-string
- \rightarrow Vortex-like solutions in field theory e.g. from spontaneously broken U(1) symmetry

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• The origins of cosmic strings:

- \rightarrow Prediction from Superstring theory: (F-) string, D-string
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- Charged complex scalar field

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Cosmic String formation (Kibble mechanism)

• Charged complex scalar field

$$V = \lambda \left(\Phi^{\dagger} \Phi - \frac{v^2}{2} \right)^2$$

- Horizon size at early time (high temperature) $d_H \propto M_p/T^2$
- we need a solution:

$$\Phi \xrightarrow{r \to \infty} \frac{v}{\sqrt{2}} e^{i\theta}$$



Cosmic String solution

• In the Abelian Higgs model

$$\mathcal{L} = D_{\mu} \Phi D^{\mu} \Phi^{\dagger} - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \lambda \left(\Phi^{\dagger} \Phi - \frac{v^2}{2} \right)^2$$



Cosmic Strings



Vilenkin and Shellard 94'

• Static string network would red-shift as

$$\rho_{\infty} \propto a^{-2}$$

• strings intercommute on collision



• overall energy density of the network scales with total energy density

$$\frac{\rho_{\infty}}{\rho_{\rm tot}} \propto G\mu$$

Cosmic String network

• We use a simplified loop distribution

$$l_i = \alpha t_i, \quad \alpha \approx 0.1$$

• After its creation each loop radiates energy at a constant rate

$$\frac{dE}{dt} = -\Gamma G \mu^2, \quad \Gamma \approx 50$$

• The loop size decreases as

$$l = \alpha t_i - \Gamma G \mu \left(t - t_i \right)$$

• The final loop density reads

$$n(l,t) = \frac{C_{\text{eff}}(t_i)}{\alpha^2 t_i^4} \frac{a^3(t_i)}{a^3(t)}$$

where the factor C_{eff} depends on the background evolution.

 \bullet the observed frequency f of GWs and time when emitting loop was created

$$f = \frac{2}{l} \frac{a(t_0)}{a(\tilde{t})} \qquad t_i = \frac{1}{\alpha + \Gamma G \mu} \left(\frac{2}{f} a(\tilde{t}) + \Gamma G \mu \tilde{t} \right)$$

• The final estimate for GW density today, reads

$$\Omega_{\rm GW}(f)h^2 = h^2 \frac{16\pi}{3f} \frac{\Gamma(G\mu)^2}{H_0^2 \alpha \left(\alpha + \Gamma G\mu\right)} \int_{t_F}^{t_0} d\tilde{t} \frac{C_{\rm eff}(t_i)}{t_i^4} \left(\frac{a(\tilde{t})}{a(t_0)}\right)^5 \left(\frac{a(t_i)}{a(\tilde{t})}\right)^3 \Theta\left(t_i - t_F\right)$$

• where the factor $C_{\rm eff}$ depends on the background evolution $H^2 \propto a^{-n}$











• GW frequency \leftrightarrow temperature

$$f_\Delta \propto rac{T_\Delta}{\sqrt{G\mu\,lpha}}$$

Detection capabilities



• slightly better numerical result

$$f_{\Delta} = (8.67 \times 10^{-9} \,\mathrm{Hz}) \, \frac{T_{\Delta}/\mathrm{GeV}}{\sqrt{\alpha \, G\mu}} \left(\frac{g_*(T_{\Delta})}{g_*(T_0)}\right)^{\frac{8}{6}} \left(\frac{g_S(T_0)}{g_S(T_{\Delta})}\right)^{-\frac{7}{6}}$$

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Extra DOF

• We add Δg_* new degrees of freedom at T_{Δ}

$$g_*(T) = \begin{cases} g_*^{\rm SM}(T) & \text{for} \quad T < T_\Delta \\ g_*^{\rm SM}(T) + \Delta g_* & \text{for} \quad T > T_\Delta \end{cases}$$

• An example with $T_{\Delta} = 200 \,\text{GeV}$ and $G\mu = 10^{-11}$



• We will model the energy budget of the universe as

$$\rho(t) = \begin{cases} \rho_{st}(t) & ; t \ge t_{\Delta} \\ \rho_{st}(t_{\Delta}) \left[\frac{a(t_{\Delta})}{a(t)} \right]^n & ; t < t_{\Delta} \end{cases}$$

- examples:
 - **()** standard radiation domination (n = 4)
 - **2** early matter domination (n = 3)
 - **③** oscillating scalar field (for non-renormalisable potential $n \to 6$) **④** ...
- experimental bounds: RD during BBN $\Rightarrow T_{\Delta} \gtrsim 5$ MeV

Non standard cosmologies



• at large frequencies

$$H^2 \propto a^{-n} \Longrightarrow \Omega_{GW}(f) \propto \begin{cases} f^{\frac{8-2n}{2-n}} & n > 10/3\\ f^{-1} & n \le 10/3 \end{cases}$$

- Cosmic strings could provide a unique and powerful tool for probing the early history of the universe.
- 2 Any departure from a flat spectrum predicted by standard cosmological evolution can be traced to the respective temperature of cosmological modification.
- This method could probe the cosmological evolution to time well before the currently available BBN data.