

# Moduli and obstructions of $\mathcal{N} = 1$ heterotic backgrounds

---

Anthony Ashmore

University of Oxford

1806.08367 w/

X. de la Ossa, R. Minasian,

C. Strickland-Constable, E. Svanes

String Pheno 2018

Calabi–Yau compactifications have large numbers of moduli

Move away from Calabi–Yau and allow non-zero **flux**

- *Most* moduli can be stabilised
- Internal spaces are **non-Kähler**

Can we say anything about **general** heterotic compactifications?

**Goal:**

Understanding of *moduli spaces*

# Heterotic string

Work in heterotic string at  $\mathcal{O}(\alpha')$

Want **Minkowski** compactifications that preserve minimal **supersymmetry**

$$M_{10} = \mathbb{R}^{1,3} \times X$$

$X$  is compact 6d space with vector bundle  $V$

- Metric  $g$
- Dilaton  $\phi$
- Gauge fields  $A$  with  $G \subseteq E_8 \times E_8$
- 3-form flux  $H$

# Hull–Strominger system

General 4d Minkowski solutions with  $\mathcal{N} = 1$  are given by the  
“Hull–Strominger system” [Strominger '86, Hull '86]

# Hull–Strominger system

General 4d Minkowski solutions with  $\mathcal{N} = 1$  are given by the “Hull–Strominger system” [Strominger '86, Hull '86]

$X$  is complex with an  $SU(3)$  structure and a conformally balanced metric

$$\begin{aligned}\omega \wedge \Omega &= 0, & \omega^3 &\propto |\Omega|^2, \\ d\Omega &= 0, & d(e^{-2\phi}\omega \wedge \omega) &= 0\end{aligned}$$

# Hull–Strominger system

General 4d Minkowski solutions with  $\mathcal{N} = 1$  are given by the “Hull–Strominger system” [Strominger '86, Hull '86]

$X$  is complex with an  $SU(3)$  structure and a conformally balanced metric

$$\begin{aligned}\omega \wedge \Omega &= 0, & \omega^3 &\propto |\Omega|^2, \\ d\Omega &= 0, & d(e^{-2\phi}\omega \wedge \omega) &= 0\end{aligned}$$

$V$  and  $TX$  are polystable holomorphic bundles

$$F_{(0,2)} = 0, \quad F \wedge \omega \wedge \omega = 0$$

# Hull–Strominger system

General 4d Minkowski solutions with  $\mathcal{N} = 1$  are given by the “Hull–Strominger system” [Strominger '86, Hull '86]

$X$  is complex with an  $SU(3)$  structure and a conformally balanced metric

$$\begin{aligned}\omega \wedge \Omega &= 0, & \omega^3 &\propto |\Omega|^2, \\ d\Omega &= 0, & d(e^{-2\phi}\omega \wedge \omega) &= 0\end{aligned}$$

$V$  and  $TX$  are polystable holomorphic bundles

$$F_{(0,2)} = 0, \quad F \wedge \omega \wedge \omega = 0$$

$H$  satisfies a Bianchi identity

$$H = i(\partial - \bar{\partial})\omega, \quad dH = \frac{\alpha'}{4}(\text{tr } F \wedge F - \text{tr } R \wedge R)$$

Difficult to find solutions! [Goldstein, Prokushkin; Fu, Yau; Becker, Sethi; Becker<sup>2</sup> et al.; . . .]

- Torsional geometries not well understood

What are the **moduli** of these solutions?

- **Deformations** of  $X$  and  $V$  that preserve SUSY
- Hermitian, complex structure and bundle moduli
- No systematic understanding until recently [Anderson, Gray, Sharp '14; Garcia-Fernandez '13; Baraglia, Hekmati '13; de la Ossa, Svanes '14]



## A holomorphic structure

Hull–Strominger system defines a holomorphic structure  $\bar{D}$  on a bundle  $\mathcal{Q}$

# A holomorphic structure

Hull–Strominger system defines a holomorphic structure  $\bar{D}$  on a bundle  $\mathcal{Q}$

- Define  $\mathcal{Q} \simeq T^{(1,0)}X \oplus \text{End } V \oplus \text{End } TX \oplus T^{*(1,0)}(X)$

# A holomorphic structure

Hull–Strominger system defines a **holomorphic structure**  $\bar{D}$  on a bundle  $\mathcal{Q}$

- Define  $\mathcal{Q} \simeq T^{(1,0)}X \oplus \text{End } V \oplus \text{End } TX \oplus T^{*(1,0)}(X)$
- Define a **differential**  $\bar{D}$  so that  $\bar{D}^2 = 0$  iff  $\bar{\partial}^2 = \bar{\partial}_A^2 = 0$  and Bianchi for  $F$  and  $H$

# A holomorphic structure

Hull–Strominger system defines a **holomorphic structure**  $\bar{D}$  on a bundle  $\mathcal{Q}$

- Define  $\mathcal{Q} \simeq T^{(1,0)}X \oplus \text{End } V \oplus \text{End } TX \oplus T^{*(1,0)}(X)$
- Define a **differential**  $\bar{D}$  so that  $\bar{D}^2 = 0$  iff  $\bar{\partial}^2 = \bar{\partial}_A^2 = 0$  and Bianchi for  $F$  and  $H$

$(\bar{D}^2 = 0)$  + polystability + conformally balanced



$(X, V)$  gives  $\mathcal{N} = 1$  solution

[Anderson, Gray, Sharp '14; de la Ossa, Svaner '14]

$H_{\bar{D}}^{(0,1)}(\mathcal{Q})$  gives the *infinitesimal* deformations

- Gives infinitesimal **massless** spectrum
- These deformations can be **obstructed** at higher orders

$H_{\bar{D}}^{(0,1)}(\mathcal{Q})$  gives the *infinitesimal* deformations

- Gives infinitesimal **massless** spectrum
- These deformations can be **obstructed** at higher orders

In the low-energy theory, the infinitesimal calculation tells you these moduli appear in the action without mass terms

The obstructions at higher orders correspond to **Yukawa couplings**

- We want to understand these higher-order contributions

# Moduli

$H_{\bar{D}}^{(0,1)}(\mathcal{Q})$  gives the *infinitesimal* deformations

- Gives infinitesimal **massless** spectrum
- These deformations can be **obstructed** at higher orders

In the low-energy theory, the infinitesimal calculation tells you these moduli appear in the action without mass terms

The obstructions at higher orders correspond to **Yukawa couplings**

- We want to understand these higher-order contributions

Analogous to complex structure defs

- Infinitesimally:  $\bar{\partial}\mu = 0$
- Higher order:  $\bar{\partial}\mu - \frac{1}{2}[\mu, \mu] = 0$

# Higher-order deformations

Higher-order deformations are difficult

- Complicated and *highly dependent* on how you parametrise the deformations



# Higher-order deformations

Higher-order deformations are difficult

- Complicated and *highly dependent* on how you parametrise the deformations

Physics guides us

- $\mathcal{N} = 1$  theory  $\Rightarrow$  4d superpotential is **holomorphic** [McOrist '16]
- Field space is **complex** with Kähler metric [Candelas et al. '15]
- Superpotential sees only **holomorphic deformations**

# The heterotic superpotential

4d heterotic theory has a GVW-like **superpotential** [Gukov et al. '99; Becker et al. '03; Cardoso et al. '03, Lukas et al. '05; McOrist '16]

$$W = \int_X (H + i d\omega) \wedge \Omega$$

# The heterotic superpotential

4d heterotic theory has a GVW-like **superpotential** [Gukov et al. '99; Becker et al. '03; Cardoso et al. '03, Lukas et al. '05; McOrist '16]

$$W = \int_X (H + i d\omega) \wedge \Omega$$

Minkowski vacuum  $\Leftrightarrow W = \delta W = 0$  on solution

- Recover  $F$ -term conditions
- $D$ -term conditions are polystability and conformal balance – *not relevant for moduli*

[de la Ossa, Hardy, Svanes '14]

(Suppress  $TX$  for now)

# Change of superpotential

Holomorphic deformations are

$$\begin{aligned}\Delta\Omega &= \iota_\mu\Omega + \frac{1}{2}\iota_\mu\iota_\mu\Omega + \frac{1}{3!}\iota_\mu\iota_\mu\iota_\mu\Omega, \\ \Delta(B + i\omega) &= x_{(1,1)} + b_{(0,2)} \\ \Delta A &= \alpha_{(0,1)}\end{aligned}$$

# Change of superpotential

Holomorphic deformations are

$$\begin{aligned}\Delta\Omega &= \iota_\mu\Omega + \frac{1}{2}\iota_\mu\iota_\mu\Omega + \frac{1}{3!}\iota_\mu\iota_\mu\iota_\mu\Omega, \\ \Delta(B + i\omega) &= x_{(1,1)} + b_{(0,2)} \\ \Delta A &= \alpha_{(0,1)}\end{aligned}$$

Generic holomorphic deformation gives

$$\begin{aligned}\Delta W &= 2 \int_X (-\iota_\mu\bar{\partial}x + \frac{1}{2}i\iota_\mu\iota_\mu\partial\omega + \dots - \frac{1}{2}\iota_\mu\partial b) \wedge \Omega \\ &+ \int_X \text{tr}(\alpha \wedge \bar{\partial}_A\alpha - 2\iota_\mu F \wedge \alpha + \frac{2}{3}\alpha \wedge \alpha \wedge \alpha + \dots) \wedge \Omega\end{aligned}$$

# Change of superpotential

Holomorphic deformations are

$$\begin{aligned}\Delta\Omega &= \iota_\mu\Omega + \frac{1}{2}\iota_\mu\iota_\mu\Omega + \frac{1}{3!}\iota_\mu\iota_\mu\iota_\mu\Omega, \\ \Delta(B + i\omega) &= \mathbf{x}_{(1,1)} + \mathbf{b}_{(0,2)} \\ \Delta A &= \alpha_{(0,1)}\end{aligned}$$

Generic holomorphic deformation gives

$$\begin{aligned}\Delta W &= 2 \int_X (-\iota_\mu \bar{\partial} \mathbf{x} + \frac{1}{2} i \iota_\mu \iota_\mu \partial \omega + \dots - \frac{1}{2} \iota_\mu \partial \mathbf{b}) \wedge \Omega \\ &\quad + \int_X \text{tr}(\alpha \wedge \bar{\partial}_A \alpha - 2 \iota_\mu F \wedge \alpha + \frac{2}{3} \alpha \wedge \alpha \wedge \alpha + \dots) \wedge \Omega\end{aligned}$$

Now want  $\Delta W = \delta\Delta W = 0$  for  $\mathcal{N} = 1$  Minkowski vacuum

- Is there some structure hiding here?

Looking for a Maurer–Cartan equation – need a **differential** and a **bracket**

Package deformation as

$$y = (x, \alpha, \mu)$$

$$y \in \Omega^{(0,1)}(\mathcal{Q}) \simeq \Omega^{(0,1)}(T^{*(1,0)}X \oplus \text{End } V \oplus T^{(1,0)}X)$$

# $\bar{D}$ and brackets

Looking for a Maurer–Cartan equation – need a **differential** and a **bracket**

Package deformation as

$$y = (x, \alpha, \mu)$$

$$y \in \Omega^{(0,1)}(\mathcal{Q}) \simeq \Omega^{(0,1)}(T^{*(1,0)}X \oplus \text{End } V \oplus T^{(1,0)}X)$$

Already have a candidate for the differential:  $\bar{D}$

$$(\bar{D}y)_a = \bar{\partial}x_a + i(\partial\omega)_{ea\bar{c}}d\bar{z}^{\bar{c}} \wedge \mu^e - \text{tr}(F_{a\bar{b}}d\bar{z}^{\bar{b}} \wedge \alpha)$$

$$(\bar{D}y)_\alpha = \bar{\partial}_A\alpha + F_{b\bar{c}}d\bar{z}^{\bar{c}} \wedge \mu^b$$

$$(\bar{D}y)^a = \bar{\partial}\mu^a$$

[Anderson–Gray–Sharp '14; de la Ossa–Svanes '14]



Appearance of  $TX \oplus T^*X$  in  $\mathcal{Q}$  suggests form of bracket

$$[y, y]_a = 2 \mu^b \wedge \partial_b x_a - \mu^b \wedge \partial_a x_b + \dots$$

$$[y, y]_\alpha = -2 \alpha \wedge \alpha + \dots$$

$$[y, y]^a = 2 \mu^b \wedge \partial_b \mu^a$$

Also have a natural pairing on sections

$$\langle y, y \rangle = 2 \mu^a \wedge x_a + \text{tr } \alpha \wedge \alpha$$

$\bar{D}$  and  $[\cdot, \cdot]$  satisfy *Leibniz identity*, and bracket satisfies *Jacobi identity* up to  $\partial$ -exact terms

# Superpotential

Change in superpotential can be written as

$$\Delta W = \int \langle y, \bar{D}y - \frac{1}{3}[y, y] - \partial b \rangle \wedge \Omega$$

# Superpotential

Change in superpotential can be written as

$$\Delta W = \int \langle y, \bar{D}y - \frac{1}{3}[y, y] - \partial b \rangle \wedge \Omega$$

$\mathcal{N} = 1$  Minkowski  $\Leftrightarrow W = \delta W = 0$  gives

$$\begin{aligned}\bar{D}y - \frac{1}{2}[y, y] - \frac{1}{2}\partial b &= 0, \\ \bar{\partial}b - \frac{1}{2}\langle y, \partial b \rangle + \frac{1}{3!}\langle y, [y, y] \rangle &= 0, \\ \partial_{i_y}\Omega &= 0\end{aligned}$$

Solutions  $(y, b)$  are **moduli**

# Superpotential

Change in superpotential can be written as

$$\Delta W = \int \langle y, \bar{D}y - \frac{1}{3}[y, y] - \partial b \rangle \wedge \Omega$$

$\mathcal{N} = 1$  Minkowski  $\Leftrightarrow W = \delta W = 0$  gives

$$\begin{aligned}\bar{D}y - \frac{1}{2}[y, y] - \frac{1}{2}\partial b &= 0, \\ \bar{\partial}b - \frac{1}{2}\langle y, \partial b \rangle + \frac{1}{3!}\langle y, [y, y] \rangle &= 0, \\ \partial_{i_y}\Omega &= 0\end{aligned}$$

Solutions  $(y, b)$  are **moduli**

- Generalisation of **holomorphic Chern–Simons** theory
- Can recast as an  $L_3$  algebra

## Summary

- Coupled moduli of the Hull–Strominger system via superpotential
- Superpotential reduces to Chern–Simons like form

## Summary

- Coupled moduli of the Hull–Strominger system via superpotential
- Superpotential reduces to Chern–Simons like form

## Still to do

- Specific examples? Can we *compute* the cohomologies?
- Are there conditions for moduli to be unobstructed?
- Quantum corrections?
- Topological theory? [Witten '91]
- New invariants? [Donaldson, Thomas '98]